

**PATENT**  
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**APPLICATION FOR UNITED STATES LETTERS PATENT**  
**for**  
**METHOD AND APPARATUS TO PREVENT SIGNAL PILE-UP**  
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The government may own rights in the present invention pursuant to grant numbers NIH-RO1-CA61880, NIH-RO1-CA58980, NIH-RO1-CA76246, and NIH-RO1-CA58980S1 from the National Institutes of Health -- National Cancer Institute.

5 This invention claims the benefit of U.S. Provisional Application Serial No. 60/045,836, by Wai-Hoi Wong, *et al.*, filed May 7, 1997 and co-pending U.S. Non-Provisional Application Serial No. 09/074,274 filed May 7, 1998. Both of those applications are hereby incorporated by reference in their entirety.

10 This application is a continuation-in-part of the co-pending Application No. 09/074,274 mentioned above, entitled "METHOD AND APPARATUS TO PREVENT SIGNAL PILE-UP" by Wai-Hoi *et al.*

## **BACKGROUND OF THE INVENTION**

### **1. Field of the Invention**

15 The present invention relates generally to the field of radiation detection and imaging technology. In particular, the present invention may be used for processing exponentially decaying pulses from a scintillation detector or a similar detector for radiation particles.

### **2. Description of Related Art**

20 When a radiation particle (gamma ray, neutron, electron, etc.) is detected in a scintillation detector, the scintillation detector will emit light, which is then converted into an electronic signal by a photosensor (*e.g.*, photomultiplier tube or photodiode). This electronic signal can then be received and processed by electronic circuits. In the period after radiation hits the scintillation detector, the scintillation light decays exponentially with a time constant  $\tau$  (the time when the light level decays to 37% of the onset level), as shown in FIG. 1.

FIG. 1 shows energy output by two gamma ray particles over time. Since the total amount of light emitted by the scintillation detector represents linearly the energy deposited by the radiation particle in the detector, the area or integral under the curves in FIG. 1 is a measure of the particle energy. As shown in FIG. 1, area 5 and area 10 define a measure of the particle energy of the gamma ray particles. Furthermore, the initial peak in the light level is also proportional to the radiation energy. Hence both the area 5 and peak V1 in FIG. 1 may be used to measure the energy of the gamma ray or radiation particle. Since the area under the curve (integral of light) includes many more light signals than the instantaneous peak light level, the integral (the total amount of light emitted) is generally used to measure the radiation energy.

As the radiation flux increases, it becomes increasingly likely that the next radiation particle may arrive at the detector while all previous events are still emitting light (FIG. 2). In this case, the identity of each individual radiation particle will be lost, and several particles will merge into one large signal, as shown in FIG. 2. In this case neither the peak level (V1 or V2 of FIG. 1), nor the integral information (area 5 or area 10 of FIG. 1) can be used to separate or measure the energy of each particle. In these situations, the detection system will fail to respond properly because of erroneous measurement.

It is known that it takes a time period of approximately  $4\tau$  to collect 98% of the scintillation light from each radiation excitation. Thus, if the next event arrives at time  $t > 4\tau$ , the pile-up-energy error on the next event will be less than 2%. Hence, to keep pile-up error small, it is desirable to minimize the chance that two events (radiation excitations) will occur in a time period less than  $4\tau$ . Since the time-lapse between two events is a random distribution (*i.e.*, the time-lapse between two events is a random variable) centered about the "average arrival time", it is generally practiced in the prior art to operate the detector so that the "average arrival time" is  $10 \times (4\tau) = 40\tau$ , to lower the random chance of having two events coming closer than  $4\tau$ . With this  $10\times$  "head-room", the probability that two events will come closer than  $4\tau$  would be approximately 10%

(using Poisson statistics). The head-room factor as a function of pile-up percentage is shown in Table 1 below:

**TABLE 1**

**HEAD-ROOM FACTOR AND PILE-UPS**

Head-Room Factor	5 times	10 times	15 times	20 times
pile-up / total (%)	18 %	10 %	6.5 %	5 %

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Thus, a 10× head-room is a reasonable choice, and is generally practiced in the prior art. When coupled with a  $4\tau$  light-collection time (system dead-time), such a prior art detector provides a measured-energy error (due to pile-up) of less than approximately 2% for approximately 90% of the time, and an energy measurement error (energy resolution) greater than 2% for approximately 10% of the time. This minimum 10× head-room ( $40\tau$ ) timing requirement means that the maximum detection-rate should be less than  $1/(40\tau)$  for the scintillation detector.

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As the role of gamma cameras expands in positron coincidence imaging, radionuclide therapy dosimetry imaging, and cardiac first-pass imaging, there is a need to significantly increase the count-rate capability of gamma cameras because of the large photon flux in these procedures. Single-crystal design, conventional detector electronics, and traditional Anger-positioning algorithm may all hinder higher count-rate imaging because of the pileup of signals of detected events in the detector. In PET cameras, especially the type using Anger-positioning detectors (all commercial systems), pileup may also be a problem.

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The present disclosure permits a scintillation detector system to operate at a much higher event-rate (count-rate) by obviating the 10× head-room factor without pile-up. The techniques of this disclosure maintain a greater event-rate with little sacrifice in the total amount of scintillation light collected, specifically at a 10 times higher radiation flux with little or no sacrifice in measurement accuracy. If the fraction of scintillation light

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plurality of computation circuits, and a plurality of sampling circuits, wherein each of the delay circuits receives a different incoming signal from a different output of a gamma camera. The delay circuit, computation circuit, and sampling circuit comprise a pile-up prevention circuit.

5           Particular embodiments will comprise a plurality of pile-up prevention circuits, and may include a digital signal processor and fast trigger connected to each of the pile-up prevention circuits. Such an embodiment may also comprise an inter-zone detection circuit connected to the fast trigger and a multi-zone-trigger processor connected to said inter-zone detection circuit, capable of centroid averaging. An exemplary embodiment  
10       may have a plurality of fast triggers.

          A method for preventing signal pile-up according to the present invention may comprise: delaying an incoming signal for a preselected time prior to passing the incoming signal; computing a weighted value of the incoming signal; and sampling the weighted value upon receipt of a triggering signal from the next radiation particle,  
15       thereby preventing signal pile-up. Computing may include amplifying the incoming signal to obtain an amplified signal, integrating the incoming signal to obtain an integrated signal, and adding the amplified signal and integrated signal to obtain the weighted value. This method thereby creates a variable signal collection time.

          Another method according to the present invention may determine position and  
20       energy information of incoming signals without pile-up. Such a method may include: delaying at least one prenormalized position signal and a total energy signal; computing a weighted value for each prenormalized position signal and the total energy signal; sampling the weighted value for each prenormalized position signal and the total energy signal, upon receipt of a triggering signal from the next radiation particle. In an  
25       exemplary embodiment, the prenormalized signals and the total energy signal may be corrected by subtracting remnant values of all previous signals.

          In yet another aspect, the present invention comprises an apparatus for dynamically detecting energy of each one of a plurality of incoming signals received

from a detector, without pile-up of previous incoming signals, including: a delay circuit connected to receive an incoming signal from the detector and to pass the incoming signal from an input to an output of the delay circuit after a time delay; a trigger circuit connected to receive the incoming signal from said detector, and for generating a triggering signal upon receipt of a subsequent incoming signals at the trigger circuit; a computation circuit connected to the output of said delay circuit for determining a weighted value of the incoming signal; a sampling circuit connected to receive the weighted value from said computation circuit, and for passing the weighted value from an input to an output of the sampling circuit upon receipt of the triggering signal; and a residual subtraction circuit connected to the output of said sampling circuit, for subtracting a residual signal value corresponding to a residual weighted value of previous incoming signals, and for providing an output signal corresponding to the energy of the incoming signal.

Another aspect of the present invention comprises an apparatus connected to a gamma camera for detecting position and energy information of each one of a plurality of incoming signals received by the gamma camera, without pile-up of previous incoming signals, including: a first delay circuit connected to receive a first incoming signal from the gamma camera, and for passing the first incoming signal from an input to an output of the first delay circuit after a first time delay; second and third delay circuit arranged like the first delay circuit to receive, delay, and pass, second and third signals; a trigger circuit connected to receive the third incoming signal from said gamma camera, and for generating a triggering signal and a timing mark upon receipt of a next third incoming signal at the trigger circuit; first, second and third computation circuits, each connected to receive an output of a respective one of the first, second, and third delay circuits, and for determining a respective weighted value for each of the first, second, and third incoming signals; first, second and third sampling circuits, each connected to receive a respective one of the first, second, and third weighted values, and for circuits passing the respective weighted value upon receipt of the triggering signal; and a digital signal processor connected to receive the first, second, and third weighted values, and for subtracting

residual signal values corresponding to residual weighted values of previous ones of the incoming signals, and for providing an output signal corresponding to a position value of the first and second incoming signals and an energy value of the third incoming signal.

5           Another aspect of the present invention resides in a method of obtaining energy information for each one of a plurality of incoming signals received from a detector, without signal pile-up, comprising the steps of: delaying an incoming signal for a preselected time; computing a weighted value of the signal after the preselected time; sampling the weighted value upon receipt of a subsequent signal; and subtracting a  
10 residual signal value from the weighted value to obtain the energy information. The residual signal value may correspond to a residual weighted value of at least one previous incoming signal, thereby preventing signal pile-up.

Yet another aspect of the present invention resides in a method of determining position and energy information of a plurality of incoming signals from a detector  
15 without pile-up, comprising the steps of: receiving a first and second prenormalized position signal and a total energy signal from the detector; delaying the first and second prenormalized position signals and total energy signal for a preselected time; computing a weighted value for each of the first and second prenormalized position signals and total energy signal after the preselected time; and sampling the weighted value for each of the  
20 first and second prenormalized position signals and total energy signal upon receipt of a subsequent first and second prenormalized position signals and total energy signal.

The methods of the present invention may be used to operate gamma-cameras (or other radiation detectors) in very high count-rate situations. The present invention  
25 includes the following features: (a) no compromise in measured energy-resolution in low count rates; (b) count recoveries and accurate energy measurement even for gamma-rays within a pile-up involving multiple gamma-rays; (c) optimal scintillation-light collection in very high count-rate situations; and (d) ability to merge with a multi-zone architecture to further increase count-rate capability. The present invention includes algorithms that



apply to all triggering gamma-rays (it is to be understood that although gamma-rays are discussed herein, the present invention applies to all types of radiation detectors), for extracting the correct energy and position of *every* triggering gamma-ray.

This disclosure also provides ways to improve energy and position resolution for the HYPER (High-Yield-Pileup-Event-Recovery) method and illustrates how to use that method in the application of coincidence imaging. Conventional single-scintillation-crystal design coupled with conventional detector electronics and the traditional Anger-position algorithm may all hinder higher count-rate imaging because of the pileup of gamma ray signals in the detector and electronics. At an interaction rate of 2 million events per second for NaI(Tl) scintillators, the fraction of nonpileup events is <20% of the total incident events. Hence, the recovery of pileup events may significantly increase the count-rate capability, increase the yield of imaging photons, and minimize image artifacts associated with pileups. The methodology of this disclosure, coined "HYPER," solves the pileup problem. One main idea of HYPER may be stated as (a) dynamic integrating: if no pileup happened, a signal from an event may be integrated for a fixed time (*e.g.*, 3 to 4 times the scintillation decay time) to fully collect all the light (signal) to maximize detection accuracy (resolution); however, if a pileup is detected, the integration of the current signal stops immediately; (b) calculating a weighted value: a weighted value is the estimated total energy that includes the energy of the current event and a residual energy from all the previous events; (c) remnant correction: calculating a pile-up free energy from two consecutive weighted values.

FIG. 20 shows an exemplary diagram of one embodiment of HYPER techniques. The energy and pre-normalized position signals from a gamma camera are delayed for a fixed period of time. The energy-signal pulse goes to a triggering circuit that generates a trigger and a timing mark at the leading edge of the event-signal. The computation circuit generates a weighted value estimating a total energy or pre-normalized position signal that includes the contribution from the current event and the remnant signals from all the previous events. The present-event weighted value is sampled and buffered, and two

consecutive weighted values are used for the correction of residual signals from all previous events.

One of the more significant discoveries concerning the HYPER method is that no matter how many previous events are piling up onto a present event, only the information from the present event and the last event are needed to subtract the pileup effects of all previous events (say, 5-6 events). This two consecutive-event method makes processing efficient and fast. The relation of the input signal, weighted value, and the residual signal is illustrated in FIG. 21.

HYPER methodology may be implemented and improved upon in several ways. Different implementations may use the same procedure and mathematical formulae as the basic HYPER method mentioned above in (a), (b) and (c). Different implementations have different advantages and trade-offs in signal noise and processing speed. When the signal noise is high, a lower noise weighted value for estimating the total energy may be extracted in several different ways:

(1) Creating a weighted-sum of an instantaneous value and an integrated value. The weighted-sum is the total energy in theory, but it can be noisy because the instantaneous value itself has more noise than the integrated value. To reduce the noise, the weighted-sum signal may go through an analog low-pass filter, and the filtered weighted-sum may be continuously or multiply sampled. Then, a digital filter such as median filter or a decay filter may be applied to those samples to generated a low noise weighted value.

(2) Weighting the integrated value. To get the total energy (present energy plus remnant energy), the integrated value may be dynamically weighted by a factor.

(3) Dynamically weighting the sum of the continuous-sampled values of the instantaneous signal. The sum signal of multiple/continuous sampled values of the instantaneous signal is the digital integration of the instantaneous signal. This method is similar to the method of (2) described immediately above, the difference being digital

integration instead of analog integration. The advantage of digital integration is that no analog discharge is necessary, which thereby eliminates analog-discharging dead-time of the system.

To further increase count-rate capability, multiple HYPER zone circuits may be implemented. Methods are presented in this disclosure to identify a zone that is hit by a real gamma. FIG. 25 shows a large detector divided into four zones (zones K, L, M and N). Since every two neighbor zones share one row or one column of photomultiplier tubes (PMT), one gamma hit may trigger two zones if its energy is collected by the shared PMTs. To find out the zone that has the "real" hit, a threshold may be set for the pre-normalized position signals of each zone to reject the "false" hit from its neighbor zone.

FIG. 26 shows an example of a two-zone design. Each zone has one independent block detector coupled to 4 PMTs. The pre-normalized position of zone 1 is  $X1=A+D$ , and the pre-normalized position of zone 2 is  $X2=C+F$ . If a gamma hits zone 2, there is no contribution to  $X1$ . However, if a gamma hits zone 1 (even in the corner crystal located in the center of PMT-B or PMT-E), to have the crystal well decoded, it still has more than 3-5% of total energy going to PMT-A and PMT-D. That means the pre-normalized position signal ( $A+D$ ) will have more than 3-5% of total energy for a real hit. Therefore, in data acquisition with a threshold  $E0$  for the total energy, if we set 3-5% of  $E0$  for a pre-normalized position signal threshold, the false hit caused by its neighbor's zone will be rejected (since it does not meet the signal threshold amount). Using a threshold corresponding to the energy's threshold for pre-normalized position is an effective way to reject false events within a zone.

The HYPER method combined with multiple zones may significantly increase the count-rate capability and dose efficiency of gamma cameras. For coincidence detection applications such as, but not limited to, positron coincidence imaging, a special additional technique is provided in this disclosure. Conventionally, integration time is fixed; therefore, the end of integration and signal-digitization (for energy and position) is

synchronized with the trigger signal (just delayed by the fixed integration time). The trigger signal is generated at the leading edge or the arrival time of the event.

Coincidence detection can be performed between the ends of signal digitization or the trigger signals. But, HYPER uses a dynamic integrating method, which means the

5 measured energy and position signals are no longer synchronized with the leading-edge of the present triggering signal because the energy and position signals are only generated or digitized at the arrival of the next event (arriving at a random time from present event).

To detect a coincidence event using the HYPER method, a trigger signal may be delayed by a fixed time and a new synchronization may be setup between the delayed trigger and the energy/position signals before they can go to a coincidence procedure.

10 The energy and position signals of all events may be buffered into a FIFO (first-in-first-out) queue. Then, a delayed trigger signal may be used to read them out individually to restore the synchronization between the events output of the FIFO. This process is shown generally in FIG. 27. The fixed delay time should be longer than the maximum  
15 integrating time. The trigger delay circuit should have good timing-resolution and a small dead time.

The HYPER methodology of this disclosure, the multiple HYPER zones method, and the event-synchronization method for coincidence detection may all improve PET camera imaging performance and facilitate gamma cameras for radionuclide therapy  
20 dosimetry imaging, cardiac first-pass imaging, positron coincidence imaging, and the simultaneous acquisition of transmission and emission data using difference isotopes with less-contamination between transmission and emission data. Those having skill in the art, with the benefit of this disclosure, will recognize additional suitable uses for these techniques.

## 25 **BRIEF DESCRIPTION OF THE DRAWINGS**

The following drawings form part of the present specification and are included to further demonstrate certain aspects of the present invention. The invention may be better

understood by reference to one or more of these drawings in combination with the detailed description of specific embodiments presented herein.

**FIG. 1** is a graphical representation of scintillator light output of gamma rays over time.

5        **FIG. 2** is a graphical representation of scintillator light output of a plurality of radiation particles having pile-up.

**FIG. 3A** is a graphical representation of exponential decay of a radiation particle over time.

10       **FIG. 3B** is a graphical representation of total integrated energy from a radiation particle.

**FIG. 4A** is a graphical representation of a weighted instantaneous and integrated signal of a radiation particle according to the present invention.

**FIG. 4B** is a graphical representation of a weighted sum of an instantaneous signal and an integrated signal of a radiation particle according to the present invention.

15       **FIG. 5A** is a block diagram of an exemplary embodiment of the present invention.

**FIG. 5B** is a block diagram of another exemplary embodiment of the present invention.

20       **FIG. 6** is a block diagram of an exemplary apparatus according to the present invention.

**FIG. 7** is an illustration of a PMT configuration that may be used in accordance with the present invention.

**FIG. 8** is a representation of boundary artifact generation.

**FIG. 9** is a block diagram of an exemplary apparatus according to the present invention.

**FIG. 10** is an illustration of study results according to the present invention.

**FIG. 11** is a graphical representation of scintillator light output of a plurality of  
5 radiation particles having continuous pile-ups.

**FIGS. 12A-12B** are graphical comparisons of energy spectra of count-rates of prior art methods (FIG. 12A) and methods according to the present invention (FIG. 12B).

**FIGS. 13A-D** are graphical representations of energy spectrum of an apparatus according to the present invention at very high count-rates.

10 **FIG. 14** is a graphical representation of energy resolution (percentage error) as a function of count-rates of prior art methods and methods according to the present invention.

**FIG. 15** is a graphical representation of detected count-rates as a function of true count-rates of prior art methods and methods according to the present invention.

15 **FIG. 16** is a graphical representation of energy spectrum of a delay line pulse clipping method.

**FIGS. 17A-F** are graphical representations of energy spectrum of a  $^{99m}\text{Tc}$  source with and without remnant subtraction.

20 **FIG. 18A** is a representation of a Monte Carlo result according to the present invention.

**FIG. 18B** is a representation of a Monte Carlo result of a pulse clipping method.

**FIG. 18C** is a representation of a Monte Carlo result of a conventional fixed integration method.

**FIG. 19A** is a flow chart of an exemplary method according to the present invention.

**FIG. 19B** is a flow chart of another exemplary method according to the present invention.

5        **FIG. 20** illustrates an exemplary circuit according to one embodiment of the present disclosure.

**FIG. 21** illustrates a pileup condition according to one embodiment of the present disclosure.

10       **FIGS. 22-24** illustrate various calculation circuits according to embodiments of the present disclosure.

**FIG. 25** illustrates a four-zone detector according to an embodiment of the present disclosure.

**FIG. 26** illustrates zone identification techniques according to an embodiment of the present disclosure.

15       **FIG. 27** illustrates a coincidence calculation circuit according to an embodiment of the present disclosure.

**FIG. 28** illustrates a high resolution time lapse measurement circuit according to an embodiment of the present disclosure.

### **DESCRIPTION OF ILLUSTRATIVE EMBODIMENTS**

20       The present invention in a broad aspect comprises a dynamic pile-up prevention technique for increasing the count-rate capability of scintillation detectors. In one embodiment, a NaI(Tl) scintillator may be used to illustrate the method. Although NaI(Tl) scintillators are discussed herein, the present invention is applicable to other detectors, such as BGO, GSO, LSO, plastic, and CsI. The dynamic pile-up technique of

the present invention has a variable detector signal collection time (deadtime), whereas conventional systems typically have a fixed deadtime of 1  $\mu$ s ( $4\tau$ ) for NaI(Tl) scintillators. The reason that 1  $\mu$ s is generally used as a fixed deadtime is that the scintillation light from a NaI(Tl) scintillator decays exponentially with a time constant of 0.24  $\mu$ s. Thus, approximately 98% of the light is collected in a 1  $\mu$ s signal-collection time. Hence, energy measurement error (energy resolution) is minimized.

From Poisson statistics, if a system is counting at a rate of  $R = 10^6$  counts per second (cps) and if the signal integration (collection) time is 1  $\mu$ s, the fraction of non-pile-up counts are only  $e^{-R\tau} = 37\%$ . Hence, approximately 2/3 of the counts are pile-up events. Though shortening the signal collection time below 1  $\mu$ s would improve the fraction of nonpile-up counts, it would also cause two problems. First, shortening the time could increase energy resolution or error (good or small energy resolution is needed to reject scatter noise), as energy resolution is inversely proportional to the square root of light collection. Second, even if the collection time is shortened electronically, scintillation light continues to be emitted by the detector from previously detected gamma ray particles. These remnant signals would be added erroneously to the signal of the following gamma ray particle.

Position sensitive detectors with Anger decoding algorithms are used in SPECT cameras and PET cameras (NaI(Tl), BGO and LSO systems) to reduce production costs. One drawback of Anger detectors is lower detection rates because all the photomultipliers (PMT's) involved in the localization of an incident event will be engaged in signal collection for a fixed time period, thereby inhibiting the detection system to process a second event incident within this time period. From energy and spatial resolution considerations, a fixed signal-collection time period of  $2\tau_s$  to  $3\tau_s$  is preferred, where  $\tau_s$  is the scintillation decay-time constant. Due to the random distribution of the time lapse between two events, the average time lapse between two events should be 10 times the signal collection time ( $20\tau_s$  -  $30\tau_s$ ) to reduce the probability of signal pile-up to 10%. For NaI(Tl) and BGO detectors, this limiting average time lapse would be 5-9  $\mu$ s, which



corresponds to a maximum count rate of 110-200 thousand-counts/sec (Kcps) per crystal or crystal-block, which has been the maximum count-rates of most gamma-cameras and BGO block detectors for many years. However, for certain applications, higher operating count-rates are preferred.

5           Recently, with the advent of coincidence positron imaging gamma-cameras, the need for increasing camera count rates has become more immediate. Currently, two approaches are used to increase gamma-camera count rates:

10           (a)     a dynamic integration method which integrates the first signal until a next event (pile-up) arrives. An estimated signal supplement is then added to the first pulse to correct for the signal collection deficit, and the same correction signal is subtracted from the next pulse to correct for the remnant light from the present pulse (Lewellen *et al.*, 1989). This method is relatively simple for a two event pile-up, but is more complicated for a multiple-event pile-up, where an additional circuit is needed for each higher multitude of pile-up.

15           (b)     a delay-line pulse-clipping (DLPC) method, which has been applied to NaI(Tl) PET cameras (Karp *et al.*, 1986). This technique has been adapted recently to coincidence imaging gamma-cameras. Generally, the DLPC technique reduces the scintillation pulse width to  $\tau$  (from  $3\tau$ ), which increases the maximum camera count rates by about 3 times. An advantage of this method is that it can be integrated into a multi-zone architecture, which splits a NaI(Tl) crystal into multiple pseudo-independent zones (as if having multiple pseudo-independent crystals) for increasing the maximum camera count rates further. It is known that the DLPC technique (Muehllehner and Karp, 1986; Karp *et al.*, 1986; Tanaka *et al.*, 1979; Miyaoka *et al.*, 1996b) may be used to alleviate the remnant signal problem, but such a technique does not remedy the degraded energy resolution problem. Furthermore, this technique only defers the onset of pile-up to a moderately high count rate.

Another proposed approach disclosed in U.S. Patent No. 5,430,406, which is hereby incorporated by reference, uses a dynamic approach to obtain a weighted sum of

the instantaneous signal and the integrated signal to depict the energy for the first event in a two event pile-up. The second event in this two-event pile-up was only used for detecting a pile-up condition to process the first event, and this second event is not processed further for recovery attempts. Hence, this method may maintain a good spectra at moderately high rates when the most of the pile-ups are two-event pile-ups, but the count-loss will be high due to the loss of the second event. Second, it can be shown theoretically that the weighted sum is generally not equal to the energy of an impinging event (See Appendix), because it also includes the remnant light from all previous events. The weighted sum is only equal to the event energy when the first event in the two-event pile-up is not a pile-up on previous events, which is a condition in contradiction with very high count-rate situations, where most events are riding on the signal of one or more previous events. Hence, the errors in this energy measurement method will be large in very high count-rate situations (when many events are part of multiple-event pile-ups), as shown in the following table derived from Poisson's statistics for NaI(Tl) for the probabilities of multiple-event pile-ups within 1  $\mu$ s:

**TABLE 2**  
**PILE-UP FRACTION AS A FUNCTION OF COUNT-RATES**

	no pile-up	2-event pile-up	3-event pile-up	4-event pile-up	5-event pile-up	6-event pile-up	7-event pile-up
1 Mcps	0.37	0.37	0.18	0.06	0.02	0	0
2 Mcps	0.14	0.27	0.27	0.18	0.09	0.04	0.01
4 Mcps	0.02	0.07	0.15	0.2	0.2	0.6	0.1

The present invention provides a new approach for preventing signal pile-up, namely, a hybrid signal processing method based on several concepts. It is known that after a scintillation crystal detects a gamma ray, the light output decays exponentially. This signal is illustrated in FIG. 3A, together with its one-standard-deviation ( $\pm\sigma$ ) error boundaries. Initially, about 40 electrons per 10 ns period are generated from a 140 KeV energy deposition (from  $^{99m}\text{Tc}$ ). The total integrated signal over a period of 1  $\mu$ s is

approximately 1000 electrons from this 140 KeV energy deposition event, as shown in FIG. 3B. For a scattered gamma ray of 70 KeV energy deposition, the initial signal would be 20 electrons per 10 ns period, and the total integrated signal would be 500 electrons. Although these numbers are specific for gamma rays detected by a NaI(Tl) scintillation detector, the same general principles, and the general curves shown in FIGS. 3A and 3B apply to other radiation particles and other radiation detectors. Thus, both the instantaneous signal and the integrated signal may be used to determine the detected gamma energy. However, the percentage-error is much less with the integrated signal collected over 1  $\mu$ s, as shown in FIG. 3B. For this reason, only the integrated signal technique is used traditionally.

The present invention includes a hybrid signal processing technique that uses both the instantaneous signal and the integrated signal together to help derive the radiation energy, since both signals contain radiation energy information.

In an exemplary embodiment, the instantaneous scintillation signal from the photomultiplier-tube (PMT) may be smoothed with a filter, for example a 10 ns time-averaging RC-filter. This smoothed instantaneous signal and the integrated signal are shown in FIGS. 3A and 3B. The smoothed instantaneous signal is amplified to the equivalent of 1000 electrons at time zero, and this amplified instantaneous signal is summed with the integrated signal (unamplified) as shown in FIG. 4A. Although the particular embodiment described herein amplifies this signal to a certain level of 1000 electrons, it is to be understood that the smoothed instantaneous signal may be amplified to various levels, depending upon the type of radiation and detector. In certain embodiments, the smoothed instantaneous signal need not even be amplified.

The resultant weighted-sum signal of the instantaneous signal and the integrated signal is a constant and is always a measure of the radiation energy if the event is not a pile-up on previous events, regardless of when the sum-signal is sampled. In an exemplary embodiment, the sampled signal amplitude may always be 1000 electrons, which is equivalent to a 140 KeV energy deposit. The weighted sum is constant in time,

because the amplified (weighted) instantaneous signal decays as  $\tau(E_0 e^{-t/\tau}/\tau) = E_0 e^{-t/\tau}$  (a property of NaI(Tl) and other scintillators), and the integrated signal increases as  $E_0(1 - e^{-t/\tau})$ , the mathematical property of integrating an exponential function. Thus, the weighted-sum is independent of the time at which the signal is sampled by the components of the present invention. Combining the instantaneous signal and the integrated signal derives the weighted sum according to the following equation.

$$\text{Weighted Sum} = E_0 e^{-t/\tau} + E_0(1 - e^{-t/\tau}) = E_0 \quad (1)$$

Given this equation for the weighted sum signal, it can be seen that at certain times, more of the signal will be derived from the instantaneous signal, while at other times, more of the signal will be derived from the exponential signal. If the weighted sum is sampled at earlier times, the variance of the measurement is higher, because more information weight is given by the instantaneous signal, as shown in FIGS. 4A-B. If the sum signal is sampled at 1  $\mu$ s, the sum converges to that of the integrated signal with a smaller variance, as shown in FIG. 4B.

In the present invention, if there is no next gamma ray detected within the conventional integration time period, for example, 1  $\mu$ s, the sum signal acquired in the duration of the present event is sampled. If a next gamma ray arrives within 1  $\mu$ s, the signal (sum) in the duration of the present gamma ray is sampled immediately before the scintillation signal of the next gamma ray reaches the signal measurement circuit. Hence, the measured signal sum relating to the present gamma ray excludes the erroneous pile-up signal from the next gamma ray. However, the signal sum may still contain signals from previous events.

A flow chart of an algorithm according to the present invention is shown at FIG. 19A. As shown in FIG. 19A, a detected signal is integrated in step 210, is amplified in step 220, and causes a trigger signal in step 230. The integrated value of the detected signal and the amplified value of the detected signal are summed in step 240. The

summed value is stored in step **250** and is sent also to the residual subtraction circuit for use in processing.

In addition, the trigger signal from step **230** is sent to determine time lapse between the present detected signal and the previous detected signal in step **260**. The sum of the previous detected signal is multiplied by an exponential of the time lapse in step **280**. This value is then subtracted from the present detected signal sum value in step **270**. The output of step **270** is the detected energy signal  $E_i$ , as shown in step **290**.

A block diagram of an exemplary embodiment of the present invention that implements this technique is shown in FIG. 5A. As shown in FIG. 5A, the exemplary pile-up prevention circuit **100** receives electrical signals from a photosensor (not shown in FIG. 5A) connected to the input of the circuit. The pile-up prevention circuit (PPC) **100** includes a timing trigger **110**, which provides a signal to the sample and discharge control circuit **170** and residual subtraction circuit **190** upon receipt of a signal. In an exemplary embodiment the timing trigger **110** may be a simple threshold discriminator or a delay-line clipping discriminator, which minimizes retriggering by the remnant signal of the same particle. Alternately, the timing trigger **110** may be another appropriate triggering device, for example, a Schmitt trigger. PPC **100** also includes a smoothing circuit **120** to smooth the received signal. In an exemplary embodiment, the smoothing circuit **120** may be a filter, such as a low-pass RC filter or a slow amplifier. Further, in certain embodiments, a smoothing circuit may not be necessary. The smoothed signal passes to a delay circuit **130**, which provides a desired delay to the signal before application to the calculation portion of the circuit. In an exemplary embodiment, the delay circuit **130** may be a delay cable or analog delay line integrated circuit. The delay circuit **130** may provide for a delay of between about 1 to 100 ns, and more preferably between about 5 to 50 ns.

As shown in FIG. 5A, the calculation portion of PPC **100** includes amplifier **140**, integrator **150**, and adder **160**. The amplifier **140** amplifies the smoothed signal, thereby providing a weighting factor to the signal. As discussed above, in an exemplary

embodiment the amplifier **140** amplifies the signal to 1000 electrons. In appropriate embodiments, an amplifier may not be necessary. The integrator **150** performs an integration of the signal over the exponential decay of the signal. The outputs of the amplifier **140** and integrator **150** are input into the adder circuit **160**, which develops the sum of these signals. The resultant sum is forwarded to sample and discharge control circuit **170**.

In an exemplary embodiment, the sample and discharge control circuit **170** may be comprised of a register or the like, and a control circuit, such as a fast analog switch, FET transistor, or the like. Upon a signal by time trigger circuit **110**, the sample and discharge control circuit **170** passes the resultant sum to analog to digital converter (ADC) **180**. The ADC **180** converts the sum to a digital signal, which is then forwarded as the weighted-sum signal to the residual subtraction circuit **190**, which subtracts any remnant signal from previous radiation events.

As shown in FIG. 5A, the residual subtraction circuit **190** includes a lapse timer **192**, which determines a  $\Delta T$  (or time difference) corresponding to the time between receipt of two detected signals. The residual subtraction circuit **190** also includes a look-up table **194** that determines an exponential value based on the time difference. This exponential value is then multiplied by the previous sum signal  $S_{i-1}$  (not shown). Finally, a subtractor **196** subtracts this remnant sum signal from the sum signal  $S_i$  relating to the present detected signal to thereby output a digital signal  $E_i$  corresponding to energy value of the present signal.

After sampling, the sample and discharge control circuit **170** discharges the output of the integrator **150** to zero immediately by sending a control signal upon receipt of the signal from the timing trigger **110** to a capacitor and switch (not shown) associated with the integrator **150**.

In operation, the timing trigger **110** senses the arrival of the initial burst of scintillation electrons from any gamma ray. Immediately, time trigger **110** sends a signal to the sample-and-discharge control circuit **170** to sample the weighted sum of the

previous gamma ray and pass the sum signal to the ADC 180. Since the timing trigger 110 has no delay, whereas the sum signal processing branch is delayed (in an exemplary embodiment, the delay may be between about 1 and 100 ns, and more preferably be between about 5 and 50 ns) with a delay circuit 130, the weighted-sum signal is sampled before the arrival of the energy signal from the next gamma ray. The delay is used to ascertain that the weighted sum of the present event is sufficiently discharged before the next event enters the integrator 150.

Hence, pile-up of the next signal onto the present is prevented, and the weighted sum will not contain the energy of the next gamma ray. The delay also provides enough time delay for the sample-and-discharge control circuit 170 to discharge (to zero) the previous integrated signal from the integrator 150 before the arrival of the next signal, which prepares the integrator 150 for the next signal. The discharge avoids pile-up of the old integrated signal onto the new integrated signal.

Since it takes a finite time for the electrical-charge output of the integrator 150 to be fully discharged to zero, especially if the previous event has high energy, this discharging time is the ultimate practical limit (processing dead-time) restricting the maximum count-rate performance of the present invention. In an exemplary embodiment, as shown in FIG. 5B, a dual-integrator design ("ping-pong") may be used as the integrator, with each integrator 150A and 150B taking turns in processing each consecutive event as follows: when a present event is detected, the circuitry will be switched via switch 145 to integrator 150A to integrate the signal of the present event; when the next event is detected, integrator 150A will stop integrating the present event and the integrator output will be sampled. At the same time, the circuits will switch the second signal to integrator 150B for integration without waiting for integrator 150A to be discharged. Hence, integrator 150A is discharged in this idle duty cycle when the integration of the next event is performed by integrator 150B. This will provide more than enough time for the integrator 150A to discharge fully.

When a third event is detected, the signal will be switched back to integrator 150A, which is fully discharged, for integration, and the integrator 150B is sampled and discharged at its leisure. Hence, with a dual-integrator “ping-pong” design, the discharging dead time is eliminated (to increase processing count-rate) and the error in energy measurement due to incomplete discharge is also eliminated.

As shown in FIG. 5B, there may be two independent channels, each of which may comprise an integrator 150A or 150B, an adder 160A or 160B, a sample-and-discharge circuit 170A or 170B, or an ADC 180A and 180B, together with an input switch 145 and an output switch 185 for selecting the processing-channel. Each channel takes turn to process the next incoming event. In other embodiments, there may be more than 2 channels present. This also provides an idle duty cycle for the integrators in each channel to be sampled and fully discharged. This dual-channel design also allows the use of a slower ADC (50% slower) to digitize the signals thus lowering cost.

The digital remnant subtraction circuit 190 may include a digital look-up-table (LUT) for determining the remnant scintillation energy (from all preceding gamma rays) that the scintillator is still emitting after the arrival of the latest gamma ray. Specifically, the remnant subtraction circuit 190 may include a register (not shown) to store a previous weighted sum, a LUT 192 for generating an exponential term, and a digital multiplication IC (not shown). The input to the LUT is the trigger-time of the present and the preceding gamma ray (actually, their time difference  $\Delta t$ ). The digital energy output of the remnant subtraction circuit 190 is thus the pile-up free energy ( $E_i$ ) of the detected gamma,

$$E_i = S_i - S_{i-1} e^{-\Delta t / \tau} \quad (2)$$

where  $S_i$  is the weighted-sum relating to the present gamma ray,  $S_{i-1}$  is the weighted sum relating to the preceding gamma ray, and  $\Delta t$  is the time interval between arrival of the present and the preceding gamma ray.  $S_i$  and  $S_{i-1}$  include the remnant signals from all preceding gamma rays. This is an exact equation, regardless of how many multiple pile-ups exist in each sum signal (see APPENDIX for the mathematical justification). The



correction term and the exactness of this algorithm, being independent of the number of multiple-pile-up remnants still being emitted by the scintillator, are significant consequences of the present invention. Thus, the scintillation detector used in connecting with the present invention permit count-rates that produce 60% multiple pile-ups, since  
5 all the remnants from multiple pile-ups can be easily corrected for by equation (2). The present invention processes all incoming events the same way, regardless of whether it is a pile-up onto a previous event or whether there is a pile-up from the next event.

All the components in a circuit according to the present invention are basic required components of conventional detector electronic systems, with the exception of  
10 the timing trigger and the residual subtraction circuit. Therefore, the real-time pile-up prevention/correction provided by the present invention is an inexpensive solution to the problems inherent in the prior art. If there is no pile-up, the system converges to a conventional integrating circuit, as the sample and discharge control circuit 170 may be set to pass the resultant sum signal at a time approximately equal to  $4\tau$  (*i.e.*,  
15 approximately  $1\ \mu\text{s}$  for the NaI(Tl) detector discussed herein). In case of a pile-up, it correctly measures the energy of all the pile-up gamma rays.

As long as the average time lapse ( $\Delta T$ ) of arrivals between two gamma rays is longer than the average detector integration time ( $T_c = 1\ \mu\text{s}$ ), the average fraction of scintillation light collected will not diminish, which means the detected energy resolution  
20 will not be compromised. Hence, a NaI(Tl) scintillation detector employing the methods and apparatus of the present invention has the potential to count up to  $1 \times 10^6$  cps without degrading the average energy resolution. A traditional detector circuit with a deadtime of  $1\ \mu\text{s}$  must lower the acceptable count-rate by 10 times, to 100,000 cps, to prevent the random arrival of the next gamma ray within  $1\ \mu\text{s}$  of the present gamma ray (Nicholson,  
25 1974). Poisson statistics shows that the non-pile-up fraction is 90% for 100,000 cps, which is why conventional NaI(Tl) detection systems cannot count over 100,000 cps. The present invention can count even faster than  $1 \times 10^6$  cps with some compromises in energy resolution.

The PPC of the present invention may be adapted for use with many different detectors and imaging systems. For example, the present invention may be used with thyroid probe, position sensitive detectors, or gamma ray detectors which detect neutrons, charged particles and gamma rays.

5           The conventional integration method (and its equivalent pulse-shaping method) has been the gold standard in the nuclear physics and nuclear medicine field for the past 40-50 years. The present invention is a significant breakthrough. This breakthrough can be applied to many nuclear detection areas, such as nuclear physics and engineering, high energy physics, nuclear medicine, industrial gauging, and oil field down-hole logging,  
10           among others. Certain applications will be discussed presently.

### **Nuclear Medicine Applications**

          The present invention has special attributes that make it particularly applicable to medical imaging, and especially gamma cameras. When a gamma ray hits a detector in a gamma camera, both energy and position information are needed for creating the image.  
15           The energy information discriminates (or qualifies) the useful imaging gamma ray from background gamma rays which have been scattered around in the body being imaged (scattering lowers the gamma energy). When a detected gamma is qualified by its detected energy, the position information locating the point of detection of the gamma ray has to be determined to create a gamma ray map (image).

20           The pile-up prevention circuit (PPC) of the present invention can determine the gamma ray energy, despite pile-ups. However, the present invention would be less useful for gamma cameras if it could not also determine the gamma ray position in pile-up situations. A particular embodiment of an apparatus according to the present method may be incorporated into a regular gamma camera to determine position information, as well  
25           as energy information in a pile-up situation.

          This embodiment may be known as a position-energy pile-up prevention (PEPP) algorithm and circuit. This solution is compatible with existing gamma camera electronics. Compatibility is important because it allows the pile-up-prevention circuit of

the present invention to be applied to present gamma cameras without major modification to their existing front-end electronics.

A regular gamma camera and its electronics generate five signals of interest for the present invention, namely  $X_+$ ,  $X_-$ ,  $Y_+$ ,  $Y_-$ , and  $Z$ .  $Z$  is the energy of the gamma ray.

- 5 The X-position and Y-position of the gamma ray can be calculated from these five signals,

$$X = (X_+ - X_-) / Z \quad (3)$$

and

$$Y = (Y_+ - Y_-) / Z \quad (4)$$

- 10 The prenormalized positions are defined as  $X' = (X_+ - X_-)$ , and  $Y' = (Y_+ - Y_-)$ . Hence, the X-position and Y-position can be calculated by a subsequent normalization with  $Z$ :

$$X = X' / Z \quad (5)$$

and

$$Y = Y' / Z \quad (6)$$

- 15 In addition to preventing energy pile-ups in scintillator-PMT's, the same PPC circuit may be applied directly to the prenormalized position signals,  $X'$  and  $Y'$ , to prevent pile-ups by generating pile-up-free  $X'$  and  $Y'$ . The mathematical proof of this property can be found in the APPENDIX. With the realization of this important property, the PPC circuit of the present invention may be connected to the prenormalized positioning  
20 signals  $X'$  and  $Y'$  in a regular camera, as well as the energy signal  $Z'$  (a fast instantaneous signal sum of all the PMT's), as shown in FIG. 6, to generate three pile-up-free signals  $X'$ ,  $Y'$ , and  $Z'$  for calculating the energy and position of the detected gamma ray despite multiple pile-ups.

The system of FIG. 6 includes a conventional gamma camera **200**, which develops prenormalized position signals  $X'$  and  $Y'$ , and energy signal  $Z'$ . Each of these signals is input into a separate pile-up prevention circuit (PPC) **210**. This PPC **210** may be the same circuit as discussed above, in connection with FIGS. 5A or 5B, or it may be a slightly modified version (as shown in FIG. 6, certain components of the circuit of FIG. 5A have been deleted for illustrative purposes). First the PPC **210** may not have a time trigger, as trigger signals for this embodiment may come directly from a fast trigger **220**, which is input to all of the PPC's **210**. This fast trigger **220** is a new addition to a conventional gamma camera. Further, the individual PPC's **210** may not have a remnant subtraction circuit, as all signals output from each PPC **210** may be input into a digital signal processor (DSP) **240**. In an exemplary embodiment, the DSP **240** may include the remnant subtraction algorithm for the  $X'$ ,  $Y'$ , and  $Z$  signals. The DSP **240** provides for faster processing. However, in certain embodiments, it may be possible to have the lookup tables within the individual PPC **210** units. The DSP **240** may also include other corrections (energy and distortion corrections) for the  $X'$ ,  $Y'$  and  $Z$  signals, and  $X$ ,  $Y$  renormalization process.

The system of FIG. 6 also includes a coincidence timer **230**, which receives a trigger signal from the fast trigger **220**. Upon receipt of the signal, the coincidence timer **230** sends a time clock signal to FIFO's **235**, which is transmitted to the DSP **240**. This time clock signal is for pile-up remnant subtraction and for coincident timing measurement in positron coincidence imaging (PET).

As discussed above in connection with FIG. 5A, the prenormalized position signals ( $X'$ ,  $Y'$ ) obtained from gamma camera **200** are delayed, amplified, and integrated within the PPC **210**. The amplified instantaneous position signals  $\tau X'(t)$  and  $\tau Y'(t)$  are summed to their integrated position-signal,  $\int X'(t)dt$ ,  $\int Y'(t)dt$ , respectively, also within the PPC **210**. These weighted sums of prenormalized positions are immediately sampled by a fast analog-to-digital converter (ADC) within PPC **210**, just before the onset of the next pile-up event. The arrival of the next gamma ray is sensed by the fast trigger **220**, which monitors the fast instantaneous total energy output,  $Z'$ . After the ADC sampling,

the integrator circuit of PPC **210** is immediately discharged so that it can start integrating the position signals coming from the next pile-up gamma ray. Alternatively, there may be two matching integrators taking turns in processing each consecutive signal to allow for an idle duty cycle for discharging the integration to zero. Further, as discussed above, there may be two matching channels of amplifiers, integrators, samplers, and ADC's to handle each successive event to allow the integrators in each channel to discharge to zero.

The fast trigger **220** also marks the arrival time of all gamma rays. The timing mark is generated by a fast clock, which in an exemplary embodiment may be a 250 MHz clock. These arrival times are sent to the DSP **240** for remnant signal subtraction and for positron coincidence detection in PET imaging. The remnants of position signals, generated by the residual light output of all preceding pile-up gamma rays, will be subtracted from the present position weighted-sum in the DSP **240** in FIG. 6 using a LUT that stores the exponential term in the following remnant subtraction operation to determine the remnant position signal to be subtracted from the  $j$ -th event:

$$RP_j = SP_{j-1} e^{-(\Delta t)/\tau} \quad (7)$$

and

$$X_{j,1}' = SP_j - RP_j \quad (8)$$

where  $SP_{j-1}$  is the last position-weighted sum and  $\Delta t$  is the time-lapse between the  $j$  and the  $j-1$  position signals. This remnant subtraction strategy automatically subtracts for all higher levels of multiple pile-ups. After the remnant subtraction, the DSP **240** calculates the normalized positioning estimation,  $X = X' / Z$  and  $Y = Y' / Z$ , where  $Z$  is the total energy after remnant correction of the present event. Alternate methods of remnant subtraction may include performing the process within the individual PPC **210** units. The DSP **240** also performs the regular correction processes for linear distortion, field nonuniformity, and regional signal-pulse-height variation, in real time. The host processor **250** is for image display and operator interface.

An apparatus according to the present invention thus decodes the position and energy of each detected event. The fast trigger 220 detects the arrival of an event, and triggers a time-mark output in the coincidence-timer 230 for the event's arrival time. This event-arrival time mark will be used to calculate the prenormalized positions and energy.

5 The calculation may be performed by the DSP 240, where exponential functions may be stored as a look-up-table. The DSP 240 also performs the energy normalization on the prenormalized positions to provide the true position signals. In this processing scheme, all the events (photopeak or true events, and scatter noise events) are included in the processing to calculate all the remnant signals from previous events, and an event's real  
10 energy is not known until after processing. Hence, the energy acceptance of an event has to be made after the decoding operation, in/after the DSP 240. If the decoded energy is higher than that of the scatter radiation noise signal, the detected radiation is accepted as a true (photopeak) event. The output of the DSP is thus the pile-up-free, energy-normalized positions (x, y) and gamma-energy (z). Since all data acquisition  
15 computers in gamma cameras accept the standard signals of (x, y, z), the proposed processing algorithm and electronic architecture is compatible with existing data acquisition computers. This compatibility is useful to adapt the scheme to existing cameras.

Implementation of the present invention in a conventional camera is very feasible,  
20 as shown in FIG. 6. A DSP circuit is standard in most cameras for processing position signals and for distortion correction. Thus, the only additional circuits needed for incorporating the methods of the present invention into a regular camera are the three channels of PPC circuits according to the present invention, a fast trigger circuit, and a clock, which increase production costs by only a few hundred dollars to achieve a 10-20  
25 times increase in count-rate capability.

Such an increase in count-rate capability would permit or improve the following imaging procedures:

5 (a) Positron imaging, because it requires the removal of a lead collimator from a gamma camera, which exposes the NaI(Tl) detector to a 10-fold or more gamma-ray flux. This forces an injected positron tracer dose to be reduced by 5-10× to prevent significant pile-up in a regular camera. With the present invention, it would not be necessary to reduce the injected dose. The positron image quality would be significantly improved as a result of this 5-10 times increase in positron counts in the image. Imaging time can also be decreased as a result of the increased count-rate capability;

10 (b) Imaging radionuclide therapy patients to deduce actual radiation dose delivered to tumors and organs so that radionuclide therapy can be improved;

(c) Dynamic first-pass cardiac imaging to study shunts, valves, the right heart and lung;

(d) The use of very short half-life tracers to reduce radiation dose and increase patient throughput;

15 (e) The use of larger NaI(Tl) detectors to reduce scanning time for whole-body imaging;

(f) Acquiring very high count-rate transmission data while the emission data is being collected. In such a scheme, a lower energy gamma source is used for the transmission data; the very high-rate transmission collection is used to minimize the effect of 'down scatter' contamination from the emission gamma. Other improvements in medical imaging would also be produced by application of the present invention.

20

The present invention also provides for further improvement of the detection rates of a gamma camera. In this embodiment, the camera signals may be split geometrically into 4 or more independent signal processing zones as if four independent cameras existed adjacent to each other. More specific details of such a system are disclosed in

25 U.S. Patent No. 5,319,204 and, U.S. Patent No. 5,453,623, the disclosures of which are

hereby incorporated by reference. The incorporation of the PEPP (position and energy pile-up prevention) algorithm of the present invention into such a multi-zone design would further increase the count-rate capability of the camera by two or three times.

Multi-zone designs have been proposed and implemented in NaI(Tl)-based PET cameras to increase maximum count-rate capability (Muehllehner *et al.*, 1995; Muehllehner and Karp, 1986; Karp *et al.*, 1986; Freifelder *et al.*, 1994). These designs depend on the fact that most of the scintillation light is distributed only to the neighboring 7-9 PMTs.

For a camera with 36 PMT's ( $6 \times 6$  configuration), the camera may be divided into four identical square zones, identified as K, L, M, N, as shown in FIG. 7. Each zone therefore has nine PMT's ( $3 \times 3$ ). Each zone may be treated as an independent camera with its own PEPP circuit according to the present invention to correct for pile-ups and image distortions. The nine PMT's in each zone may be grouped into one signal-triggering line. With four zones, there would be four signal-triggering lines,  $S_K$ ,  $S_L$ ,  $S_M$ ,  $S_N$ . If the first gamma ray is detected in zone K, the PEPP circuit of zone K would be turned on to measure its position and energy. If a second gamma is detected in zone K within the pile-up time, the PEPP circuit would correctly measure the position and energy of both gamma rays.

If the second gamma ray strikes zone L while zone K is processing a prior count, two scenarios may occur. These two scenarios are addressed in detail:

(1) In the first scenario, the gamma rays detected in zone K and zone L are both far from the zone boundaries between K and L, such that there is little light spilled from one zone to the next. Hence, there is little or no signal interaction between the two events to cause positioning errors. Both gamma rays may be processed independently by the PEPP circuit of each zone to obtain the correct energy and position.

(2) In the second scenario, the first gamma ray detected in K is so near the K-L boundary that the PEPP circuits in both zone K and zone L are triggered, as shown in



FIG. 8. In this case, zone K may decode a valid position, but zone L would decode an artifact-count at the K-L boundary near the event in K, as shown in FIG. 8. When the second gamma hits zone L later, at location A in FIG. 8, the remnant signal of the fictitious count created in the L-PEPP circuit would be subtracted from the signal of the second gamma ray (A), so that there is no interference in determining the position of the second gamma ray. Thus, in this scenario, both the first (K) and second (L) gamma rays would be correctly measured, but an additional artifact count will be created at the K-L boundary.

With these inter-zone light-spill considerations, the processing circuitry for the multi-zone pile-up prevention (MZPP) camera is shown in FIG. 9. As shown in FIG. 9, each zone may have its own PEPP circuit 310. In addition, each zone may have its own DSP 315. These four PEPP circuits may be monitored by an inter-zone coincident detection circuit 330 and a multi-trigger processor 340. The interzone coincident detection circuit receives input from the fast timing triggers of each of the zone PEPP circuits 310. The multi-trigger processor 340 receives zone position signals from each of the zone PEPP circuits 310. In an exemplary embodiment, the processor 340 may be a DSP. The output of processor 340 is input into a main DSP processor 350, which provides information to image buffer 360, interface 370 and Host CPU 380. The main DSP processor 350 serves two functions: (1) it merges the four zone-images into one image, and (2) it performs a final distortion correction for the combined image.

To eliminate the problem of additional artifact counts at a border, the interzone coincidence-triggering circuit 330 (between any two or more zones) may be used to detect the simultaneous triggering of two or more zones caused by an event detected near the boundary. When two or more neighboring zones are triggered by a single event near the boundary, the PEPP circuit 310 in each zone would respond (as originally designed) as if an independent event is detected in each zone by generating their own position signals,

$$X_K = (X_{K+} - X_{K-}) / Z_K \quad (8)$$

and

$$X_L = (X_{L+} - X_{L-}) / Z_L \quad (9)$$

Furthermore, the simultaneous triggering of both zones would also activate the inter-zone coincidence circuit **330** which would then feed these two or more independent position signals into multi-zone-trigger processor **340**, which performs a centroid averaging for these two or more position signals. Mathematically, this averaging is equivalent to combining the two or more zones and using the regular Anger positioning method over this larger domain, as shown in the following equation:

$$X = (X_K Z_K + X_L Z_L) / (Z_K + Z_L) = \{ (X_{K+} - X_{K-}) + (X_{L+} - X_{L-}) \} / (Z_K + Z_L) \quad (10)$$

This averaged position is stored in the main processor **350** that stores the composite camera image (four zones combined), whereas the individual zone position signals ( $X_K$  and  $X_L$ ) are discarded. This method thereby eliminates the boundary-artifact event and also provides a better estimation of the position, since all of the scintillation light emitted would be used for computing the position (including light spilled into the adjacent zone).

Inter-zone light spill also exists in the two high count-rate cameras using a multi-zone design because, as long as a single NaI(Tl) crystal is used, light will be distributed from one zone to the next. Technical solutions are readily achievable as indicated by ADAC MCD and UGM SPECT and PET cameras (Muehllehner *et al.*, 1995; Glass *et al.*, 1996; Freifelder *et al.*, 1994). The present invention, using a monitor to detect coincidence triggering of two zones is simple, straightforward to implement, and usable with the pile-up prevention technology of the present invention.

An increase in detection rates by 2-3 times is expected over the single-zone implementation, if the distribution of the scintillation light is mainly limited to a circle 4.5" in diameter centered on the point of detection. When this light-spread condition is met, only the gamma rays hitting the shaded area shown in FIG. 10 can trigger or affect

both zones. Hence, the probability that a gamma ray will hit the independent area is 21/36, while the probability that it will hit the light-spill-over area is 15/36. Thus, the count-rate enhancement obtained by splitting the camera signals into 4 zones is (21/36) ×4 or 2.3 times.

Shown in FIG. 19B is a flow chart for the method of the present invention relating to this multi-zone pile-up prevention embodiment. As shown in FIG. 19B, an instantaneous detected signal is provided to four independent PEPP's in step 410. Then it is determined whether this signal creates a multi-trigger event in step 420. If it is not a multi-trigger event, control passes to step 430, in which it is determined whether it is a single-trigger event. If it is a single-trigger event, control passes to step 440 and single active zone position processing is performed, as discussed above.

If it is determined that it is a multi-trigger event, in step 450 active zones are combined into a larger domain for processing. Then, the large domain position processing is performed in step 460, as discussed above. The result of either path of this method results in a determination of the true and accurate position of the instantaneous signal in step 470.

Combining the pile-up prevention method and the multi-zone method, the count-rate of a camera according to the present invention may reach 4,000,000-5,000,000 cps. With optimization and further development of the PEPP circuit as discussed earlier, the count-rate may be extended further. Monte Carlo simulation indicates that  $2 \times 10^6$  cps is achievable with the pile-up prevention method alone, and a combined count-rate of  $2.3 \times 2 \times 10^6 = 4.6 \times 10^6$  cps may be achievable, a significant improvement over the  $10^5$  cps achievable in a regular camera.

The present application of scintillation detectors can be divided into four main areas:

- (i) industrial gauging, including thickness, level and density gauging;

(ii) surveying, including oil field logging, nuclear reactor monitoring, nuclear fuel cycle monitoring and airport luggage inspection for explosives;

(iii) research applications in nuclear physics, high-energy physics, medicine, and industry; and

5 (iv) medical uses (gamma cameras, PET cameras, bone scanners, thyroid probes, general purpose probes, monitors, and dosimeters).

A new and potentially important application of the present invention is in airport surveillance to look for plastic explosives in luggage using, for example, neutron activation techniques. The ability to detect radiation 10-20 times faster would allow (a)  
10 detectors to be used in a much higher radiation areas to extend the usefulness of the detector system, (b) the ability to count faster would increase the speed of data acquisition to shorten the data collection time, *e.g.*, in an airport bomb surveillance system, a 10-20× stronger neutron source can be used to scan luggage faster, so that luggage scanning times can be reduced by 10-20 times to improve throughput.

15 In industrial gauging application, the scintillation detector is used in conjunction with radiation sources to measure the quantities or density of materials being processed. For example, while molten cement or molten metal is poured into a container, a nuclear level gauge may be used to stop the pouring process when a specific level is reached. Another example is "thickness gauges" using gamma-ray backscatter detectors in  
20 thin-film processing in the plastic and paper manufacturing industry (especially for high-value films such as video-tapes and electronic capacitor films). Quite often, the film to be scanned is manufactured in large-area sheets, and the detectors together with the radiation source must scan the entire sheet. The detector's ability to detect gamma rays at very high count-rate (10-20×) means that a very intense gamma-ray source (10-20×) can  
25 be used, which in turns translates to a 10-20× faster data acquisition. Thus, the time spent on the inspection-scanning processes or manufacture-control processes can be reduced by 90%, thereby increasing production rate.

Scintillation detectors are also widely used in radiation monitoring in nuclear reactors, and in oil field borehole logging to survey rock/hydrocarbon structures along the borehole in oil exploration. A faster counting system would increase the information collected or decrease the data collection time, which would lower the cost of data acquisition (oil rig time is very expensive).

Scintillation detectors are also widely used for nuclear applications, for example: (a) reactor monitoring of liquid and gaseous streams to look for isotopes, (b) fuel rod cladding failure, (c) isotope scanning for irradiation fuel to determine power distribution and migration of fission products in the reactor core, (d) reactor fuel fabrication and quality control, (e) spent fuel reprocessing, and (f) management of nuclear waste. Since these nuclear-reactor related detectors are used in very high radiation flux areas, this high count-rate invention is potentially very useful in the nuclear energy business.

In research laboratories, a faster counting detector system is always welcome, as it allows new experiments to be performed. Nuclear-electronics instrumentation companies that market detector electronic modules for research laboratories may wish to add this electronic invention to their catalog.

In medical imaging applications, present gamma cameras are adequate for imaging diagnostic quantities of single photon tracers, but their limited count-rate capabilities introduce image artifacts and count-loss, and degrade image quality when gamma ray flux is high. These situations may occur in:

(a) positron imaging, which exposes the NaI(Tl) detector to a 10-fold or more gamma-ray flux, with the removal of the lead collimator from a gamma camera for positron imaging;

(b) imaging radionuclide therapy patients to deduce the actual radiation dose delivered to tumors and organs for improving radionuclide therapy treatment;

(c) dynamic first-pass cardiac imaging to study shunts, valves, the right heart and lung;

(d) the use of very short half-life tracers to reduce radiation dose and increase patient throughput;

5 (e) the use of larger NaI(Tl) detectors to reduce scanning time for whole-body imaging; and

(f) acquiring very high count-rate transmission data while the emission data is being collected (the very high-rate transmission collection minimizes the effect of 'down scatter' contamination from the emission gamma).

10

With this invention, the count-rate capability of gamma cameras can be enhanced by 20-40 times, which would allow (or improve upon) the above useful imaging protocols. The new capability may reduce reliance on the use of PET costing \$2,600,000 to purchase and \$250,000/yr to operate. Another application of the present invention is for the low-cost, high-resolution PET detection system of PMT-quadrant sharing design, as shown in U.S. Patent Nos. 5,319,204 and 5,453,623, to compensate for its lower count-rate.

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Thus, the present invention will open up new and exciting clinical applications for gamma cameras, which is especially important in today's healthcare environment when clinics must perform more with less money.

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### **HYPER Techniques and Applications**

As the role of gamma cameras expands in positron coincidence imaging, radionuclide therapy dosimetry imaging, and cardiac first-pass imaging, there is a need to significantly increase the count-rate capability of gamma cameras because of the large photon flux in these procedures. Conventional detector electronics and the traditional Anger-positioning algorithm, however, hinder higher count-rate imaging because of the

25

5 pileup of signals of the detected events in the detector. Multiple events may be merged into one artifact event in the detection system; hence, true events may be lost while artifact events are generated. The pileup problem also exists in dedicated PET cameras, albeit to a lesser extent. This disclosure allows one to recover a pileup-free pulse signal even in multiple-pileup conditions.

10 This disclosure provides several ways to implement the formulae developed in the HYPER method, thereby either improving the energy and position resolution or increasing the processing speed for pile-up prevention techniques. This disclosure also introduces a zone-identification methods for the application of multiple HYPER circuits to further increase the count-rate in gamma and PET cameras. It also explains how to use the HYPER method in the application of coincidence imaging.

15 The HYPER method involves a set of formulae that solve the signal pile-up problem in scintillation detectors for both energy and position signals despite the pileup of multiple signals. The spirit of the methodology is as follows: (a) dynamic integration – the signal integration triggered by a present-event stops immediately before the next event is detected; (b) calculating a weighted value for estimating a total energy in the scintillation, which includes the energy of the current event and a residual energy from all the previous events; and (c) remnant correction is used to calculate a pile-up free energy from two consecutive weighted values.

20 The above three steps may be used to create a pile-up free signal. This disclosure explains several ways to implement and to improve the measured resolution of the weighted value for estimating the total energy by: (1) applying an analog filter to the weighted-sum of an instantaneous value and an integrated value; the weighted-sum signal is multiple sampled, and another digital filter may be applied to those samples to get a better weighted value; (2) dynamic, digital weighting the integrated value; (3) dynamic weighting the sum of the multiple-sampling (repeated sampling) of the values of the instantaneous signal.

In certain embodiments of this disclosure, a dynamic integrating method is used, which means that the measured energy and position signals of the present-event are no longer synchronized with the trigger. This is the case because the energy and position signals of the present-event are measured at the arrival-time of the next-event (at a random time), and the trigger is issued at the arrival time of the present-event. To detect a coincidence event using the HYPER method, the original trigger signal may be delayed, and a new synchronization may be setup between the trigger and the energy or position signals before they go to a coincidence procedure. The improved HYPER method, multiple HYPER zones method, and this synchronization method may all facilitate the use of PET for applications such as, but not limited to, positron imaging, gamma cameras for radionuclide therapy, dosimetry imaging, cardiac first-pass imaging, positron coincidence imaging, and the simultaneous acquisition of transmission and emission data using difference isotopes with less-contamination between transmission and emission data.

## **Theory**

Every exponentially decreasing pulse signal, for example from a scintillation detector, includes a present pulse signal and a remnant signal from all the previous pulse signals. With reference to FIG. 2, the total light signal after event 4 is the sum of the light signal of event 4 and a remnant light signal.

$$\frac{S_n}{\tau} e^{-(t-t_n)/\tau} = \frac{E_n}{\tau} e^{-(t-t_n)/\tau} + \frac{R_n}{\tau} e^{-(t-t_n)/\tau}$$

where  $S_n$  is the total energy inside the detector after event-n without post-pileups,  $E_n$  is the energy of the event n,  $R_n$  is the remnant energy from all the previous events, and  $\tau$  is the decay time. So, we have  $S_n = E_n + R_n$ .



By the time of  $t_{n+1}$ , the pulse signal becomes  $\frac{S_n}{\tau} e^{-(t-t_n)/\tau} = \frac{S_n}{\tau} e^{-(t_{n+1}-t_n)/\tau} e^{-(t-t_{n+1})/\tau}$ .

This pulse signal will become the remnant signal to event (n+1). That means

$$\frac{R_{n+1}}{\tau} e^{-(t-t_{n+1})/\tau} = \frac{S_n}{\tau} e^{-(t_{n+1}-t_n)/\tau} e^{-(t-t_{n+1})/\tau}. \text{ Thus, we have } R_{n+1} = S_n e^{-(t_{n+1}-t_n)/\tau}.$$

Therefore, we have  $S_{n+1} = E_{n+1} + R_{n+1} = E_{n+1} + S_n e^{-(t_{n+1}-t_n)/\tau}$ . So, the energy of the event (n+1) can be calculated from the following equation:

$$E_{n+1} = S_{n+1} - S_n e^{-(t_{n+1}-t_n)/\tau} \quad (2.1)$$

This method is good for every exponentially decreasing pulse signals with multiple pileups. It also applies to the pre-normalized position signals in Anger positioning gamma cameras, since the pre-normalized position signals are exponentially decreasing too.

$$PX_{n+1} = SX_{n+1} - SX_n e^{-(t_{n+1}-t_n)/\tau} \quad (2.2)$$

$$PY_{n+1} = SY_{n+1} - SY_n e^{-(t_{n+1}-t_n)/\tau} \quad (2.3)$$

where  $PX_{n+1}$ ,  $PY_{n+1}$  are the pileup-free pre-normalized position signals for event (n+1) and  $SX_{n+1}$ ,  $SY_{n+1}$  are the total estimated pre-normalized position signals.

### Circuit design

The present disclosure, in a broad aspect, comprises a dynamic pileup prevention technique for increasing the count-rate capability of Anger type cameras with scintillation detectors such as, but not limited to, NaI(Tl), BGO, GSO, LSO, plastic, and CsI. The dynamic pileup prevention technique has a variable detector signal collection time (deadtime), whereas conventional systems typically have a fixed deadtime ( $4\tau$ ,  $1\mu s$  for NaI scintillator). The integration of the first event stops immediately if the next event is detected

with a certain period of time ( $4\tau$ ). FIG. 20 shows an exemplary diagram of an appropriate circuit.

### Computation circuit

5 The computation circuit in FIG. 20 calculates a total value of a first and second incoming pre-normalization signal and a third incoming energy signal.

$$S_n = \int_{t_n}^{\infty} \frac{S_n}{\tau} e^{-(t-t_n)/\tau} dt = \int_{t_n}^t \frac{S_n}{\tau} e^{-(t-t_n)/\tau} dt + \tau \frac{S_n}{\tau} e^{-(t-t_n)/\tau} \quad (2.4)$$

$$S_n = \int_{t_n}^{\infty} \frac{S_n}{\tau} e^{-(t-t_n)/\tau} dt = \frac{1}{1 - e^{-(t_{n+1}-t_n)/\tau}} \int_{t_n}^{t_{n+1}} \frac{S_n}{\tau} e^{-(t-t_n)/\tau} dt \quad (2.5)$$

### Weighted-sum signal

10 Equation (2.4) indicates that the weighted-sum signal of the integrated signal and the instantaneous signal amplified by  $\tau$  is a constant. The weighted-sum signal is always a measure of the total radiation energy before the next event is detected, regardless of when the sum-signal is sampled. Since the instantaneous signal is more noisy than the integrated signal, a low-pass smoothing circuit may be applied to the weighted-sum signal to reduce the noise. To further increase the signal to noise ratio, the weighted-sum signal may be multiple/continuously sampled, and a digital filter may be applied to the samples to get a better estimated value.

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In one embodiment, the filter may be a median value filter or a weighting output of these samples. For example,  $S_n = W_1 S_n^1 + W_2 S_n^2 + \dots + W_m S_n^m$ , where  $S_n^i$  indicates the i-th sample of the weighted-sum signal  $S_n$  may be used. The sum of the weighting factor is ( $W_1 + W_2 + \dots + W_m = 1$ ). Since a later sample of the weighted-sum has better statistics than an earlier sample, the weighting factor for an earlier sample should be smaller than a later sample. For example, with  $W_i = W_{i+1} e^{-T/\tau}$  (where  $T$  is the time lapse between samplings), the weighting factor follows the same exponentially decay as the pulse signal.

20

FIG. 24 shows a diagram of an appropriate decay filter. The number of samples (N) is not fixed. Thus, to obtain the estimated total energy, the accumulated output signal may be normalized by:

$$\frac{1 - e^{-N \Delta T / \tau}}{1 - e^{-\Delta T / \tau}}$$

5 In one embodiment that leads to a a simple electronics design, one may buffer the several latest samples of the weighted-sum signals and create an average of the buffered signals for the estimated total energy.

### **Weighted integrated-value**

10 Equation (2.5) indicates that the total estimated energy may be calculated by weighting the integrated value (integrating the pulse signal from  $t_n$  to  $t_{n+1}$ ). The weighting factor may be defined by the time lapse between the two events n and n+1. One possible disadvantage of this method is that the system needs good timing resolution because the weighting factor is sensitive to time resolution — especially with a small time lapse. Additionally, it may need a real-time calculation or a look-up-table to generate the result.

### **Weighted digital integration**

15 Digital integration means the instantaneous signal is continuously sampled, and the samples of each event are accumulated in digital form. Unlike analog integration, there is no dead-time of integration discharging for digital integration. However, this method requires a fast analog to digital convertor and a digital accumulator. The rest of the  
20 calculation for estimated total energy is the same as described above — the accumulated value may be weighted by a factor that is defined by a time lapse of two continuous events.

### **Remnant correction**

25 Remnant correction is an important step in creating a pileup-free signal, especially for multiple/continuous pileups. In remnant correction, the residual signal is subtracted

from the total estimated energy. The residual term may not be a small correction, and without the residual-subtraction, the weighted value may be larger than the real energy of the present event by two times or more since the preceding pulse can itself be riding on several other earlier signals, which makes the proceeding weighted value much larger than the energy of the preceding signal and the energy of the present signal.

FIG. 22 illustrates one embodiment of remnant correction, where  $T_{n+1}$  is the time lapse between event (n+2) and event (n+1), ( $T_{n+1} = t_{n+2} - t_{n+1}$ ). A remnant signal may be created by weighting the  $S_{n+1}$  with a factor  $\exp(-T_{n+1}/\tau)$ , which may be given by a look-up-table (LUT) of  $T_{n+1}$ . This remnant signal may be stored in a register for the residual correction of the next event.

If the total energy (weighted value)  $S_n$  is generated from the integrated value  $I_n$ ,

$$S_n = \frac{1}{1 - e^{-(t_{n+1}-t_n)/\tau}} I_n, \quad \text{where} \quad I_n = \int_{t_n}^{t_{n+1}} \frac{S_n}{\tau} e^{-(t-t_n)/\tau} dt,$$

From equation (2.1), we have

$$E_{n+1} = \frac{1}{1 - e^{-T_{n+1}/\tau}} I_{n+1} - \frac{e^{-T_n/\tau}}{1 - e^{-T_n/\tau}} I_n = \frac{1}{1 - e^{-T_{n+1}/\tau}} I_{n+1} - \left( \frac{1}{1 - e^{-T_n/\tau}} I_n - I_n \right), \text{ where}$$

$$T_n = t_{n+1} - t_n.$$

FIG. 23 illustrates this. The look-up-table (LUT) gives the weighting factor of  $\frac{1}{1 - e^{-T_n/\tau}}$  for the input of  $T_n$ . All the remnant signals (from multiple pileups) that goes to the next event may be generated by subtracting the integrated value from the total estimated value. This remnant signal may be stored in a register and used for the residual correction of the next event.

#### For signals with two or more decay time

Some detectors emit signals that follow two or more exponential decay-time constants. In these instances, the extractions of pileup-free energy and position may be

more complicated. The same general methodology of the common mono-exponential case as described above, however, may be used. Each integration may be sampled twice (or more) over the course of integration. The additional sampling provides another set of equations to derive the relative magnitudes of the two (or more) exponentials in the total energy (present event plus remnants). Even though for an isolated event (non-pileup), the relative magnitudes for the event are constant and known, the remnant energy in a pileup case will have different relative magnitudes for the two exponentials depending on the time-lapse between consecutive events. The additional samplings allow one of ordinary skill in the art to derive the relative strengths of the exponential.

### High resolution time measurement

High resolution time measurement may be an important issue for pileup processing. One needs to measure the time lapse between every two continuous events for the remnant corrections of both energy and position signals: see equations (2.1), (2.2) and (2.3). And, one also needs the time lapse information to calculate a total energy signal inside the scintillation if the total value is calculated from the integrated value by equation (2.5).

Suppose the time lapse measurement error is  $\delta$ . Then, the measured remnant correction value will be  $S_n * e^{-(t_{n+1}-t_n+\delta)/\tau} = e^{-\delta/\tau} * (S_n * e^{-(t_{n+1}-t_n)})$ . It is  $e^{-\delta/\tau}$  times of the real remnant value. If  $\delta = 10\text{ns}$  and  $\tau = 200\text{ns}$ , then one has  $e^{-\delta/\tau} = e^{-10/200} = 0.95$ . That means the measured remnant value is 5% less than the real remnant value, and this 5% remnant value still remains in the corrected output result. In a high count-rate situation, since the average time lapse between two events is small and the remnant value is big, the 5% error of remnant value becomes significant.

This disclosure involves methodology to measure the time lapse with a resolution, in one embodiment, less than 1 ns. This means that the remnant correction error may be less than 0.5% for  $\tau = 200\text{ns}$ . A counter with a regular resolution (for example 12ns) combined with a high resolution delay line (for example, 1ns resolution) may be used to measure the time lapse in high resolution. One could use a 1ns clock to count the time

directly, but it would need faster ECL circuits, which are expensive and may increase digital noise level. A better method is described below.

An embodiment of this disclosure is to use a slow clock to measure the time lapse to get a raw time  $T_m$ . In FIG. 28,  $T_m$  is 7 cycles of the slow clock. The rising edge of each event changes the status of a 'flag' signal (from 'high' to 'low', or from 'low' to 'high'). This flag signal inputs into a delay line with multiple tap outputs, which has, for example, a 1 ns time delay difference between two continuous taps. The rising edge of the first clock after each event-trigger latches the delay line outputs. To measure the time  $Td_n$  between the rising edge of the 'flag' signal and the first clock signal, one may count the number of delay line taps with 'high' output. This number multiplied by the tap delay step (1ns) is the time  $Td_n$  (10ns in FIG. 28).

The same methodology may be used to measure the time  $Td_{n+1}$  between the falling edge of the 'flag' signal and the following first clock — one may count the number of delay line taps with 'low' output and multiply by the tap delay step. In FIG. 28, this is 3ns. The time lapse  $T$ , or  $(t_{n+1} - t_n)$ , between event (n) and event (n+1) may be calculated by  $T = T_m + Td_n - Td_{n+1}$ . In this way, the measured  $T$  has the same timing resolution as the delay line's, which is a 1ns resolution.

In general, this methodology involves determining an approximate time between the two consecutive events using a first clock. Then, a first and second correction to the approximate time is determined using a second clock, the second clock having a finer time resolution than the first clock. Finally, the time between two consecutive events is found by applying the first and second corrections to the approximate time.

With the benefit of this disclosure, those having skill in the art will recognize that this methodology will apply to clocks having different clock resolutions (*i.e.*, other than 1 ns example given above).

## Ultra High Count-rate Application

To further increase the count-rate in a large detector with many photomultiplier tubes (PMTs), multiple HYPER zone circuits may be implemented. A detector may be divided into a plurality of smaller HYPER zones by grouping PMTs into multiple smaller Anger cameras in electronics. Some PMTs may be shared by two or more HYPER zones. Events collected by shared PMTs may trigger two or more HYPER circuits at the same time.

This disclosure provides a method to identify a zone that is hit by a real gamma and reject false hits within other zones. FIG. 25 shows a detector divided into four zones (zones K, L, M and N). Since every two neighbor zones share one row or one column of PMTs, one gamma hit may trigger two zones if its energy is collected by the shared PMTs. To find out the zone that has the “real” hit, a threshold may be set for the pre-normalized position signals of each zone to reject the false hit from its neighbor zone.

FIG. 26 shows an example of a two-zone design. Each zone has one independent block detector coupled to 4 PMTs. The pre-normalized position of zone1 is  $X1=A+D$ , and the pre-normalized position of zone2 is  $X2=C+F$ . If a gamma hits zone2, there is no contribution of light to  $X1$ . However, if a gamma hits zone1 (even in the corner crystal located in the center of PMT-B or PMT-E), it may still have, in one embodiment, more than about 3-5% of total energy going to PMT-A and PMT-D. That means that the pre-normalized position signal ( $A+D$ ) will have more than 3-6% of total energy of a real hit. Therefore, in data acquisition, a “real” hit in a zone should pass both the energy-acceptance threshold and pre-normalized position-acceptance threshold, which in this example would be about 3-6% of  $E0$  (the zonal energy-acceptance threshold). However, a false hit caused by its neighbor’s zone may be accepted by the energy threshold but it will not be accepted by the pre-normalized position threshold that is set to 3-6% of  $E0$ .

The simultaneous acceptance of the energy threshold and the pre-normalized position threshold is one way to reject the false trigger (event) within a zone. This technique leads to very robust discrimination because it is immune to differences in

electronic-gain balance between different zones. This design also reduces electronic complexity and cost.

### **Coincidence Application**

The HYPER method combined with multiple zones may significantly increase the count-rate capability and dose efficiency of PET and gamma cameras. But for coincidence detection applications and other timing-detection applications such as positron coincidence imaging, a special additional technique is provided. In conventional methods, the signal integration time for each event is fixed. Therefore, the end of integration and signal digitization (for energy and position) is synchronized with the trigger signal (delayed by the fixed integration time). The trigger signal is generated at the leading edge or the arrival time of the event. Coincidence detection may be performed between the end of digitization signals or the trigger signals.

HYPER, however, uses a dynamic integrating method, which means that the measured energy and position signals are no longer synchronized with the leading-edge of the present triggering signal because the energy and position signals are generated or digitized at the arrival of the next event (arriving at a random time from present event). To detect a coincidence event using the HYPER method, the original trigger signal may be delayed by a fixed time, and a new synchronization process may be setup between the delayed trigger and the energy/position signals before they go to a coincidence procedure.

In one embodiment, the synchronization process works as follows. The energy and position signals of all events may first be buffered into a FIFO (first-in-first-out) temporary memory. Then, a delayed trigger signal may be used to read them out individually. Since the readout of the energy and position signals from the FIFO queue is time-latched by the delayed trigger, and since the delayed trigger signal is synchronized with the trigger signal (except for the fixed delay-time shift), the energy and position signals may be synchronized to the trigger signal. This is shown in FIG. 27.



The fixed delay time should be longer than the maximum integration time. The trigger-delay circuit should have a good time resolution and a small dead time since the delayed trigger signal is synchronized with the trigger signal (except for the fixed delay-time shift).

- 5           The following table gives an example of processing five continuous events by using a FIFO and the delay technique of this disclosure to restore the synchronization between the event trigger signal and the data output signal (except for the fixed delay-time shift).

<b>Event number</b>	<b>1</b>	<b>2</b>	<b>3</b>	<b>4</b>	<b>5</b>
time of event arrival	0	0.2	0.7	1.5	2.0
event integration (sampling) time = input time of event into FIFO	0.2	0.7	1.5	2.0	3.0
Event readout time from FIFO	1.5	1.7	2.2	3.0	3.5
Waiting (holding) time of event in FIFO	1.3	1.0	0.7	1.0	0.5

- 10           In the above table, the units for time are microseconds. The event integration (sampling) time is the next event's arrival time if the current event is piled-up by the next event. If there is no pileup on the current event, it will be integrated for a fixed maximum time period. In this example, 1 microsecond was used. The waiting time is the time difference between the FIFO readout time and the event input time.

- 15           The order of the data in a FIFO is the same as the order of events arriving. Hence, the linear structure of the FIFO is used to line up all the events in the order that they arrive. Then, a large fixed delay (larger than the longest integration time) is applied for

the event-trigger, and the delayed event trigger is used to “clock” the readout of events from the FIFO. Hence, each event in the FIFO may be clocked out (readout) synchronously with the event triggering time (exact for a fixed delay). It is the waiting (holding) time of an individual event in the FIFO (as controlled by the delayed trigger) that synchronizes the FIFO readout time to the original event-arriving time.

The following examples are included to demonstrate preferred embodiments of the invention. It should be appreciated by those of skill in the art that the techniques disclosed in the examples which follow represent techniques discovered by the inventor to function well in the practice of the invention, and thus can be considered to constitute preferred modes for its practice. However, those of skill in the art should, in light of the present disclosure, appreciate that many changes can be made in the specific embodiments which are disclosed and still obtain a like or similar result without departing from the spirit and scope of the invention.

#### WORKING EXAMPLES

In studies with an exemplary embodiment, the present invention was used in connection with a 3”x 4” NaI(Tl) scintillator with a photomultiplier from Ortec Corporation (Oak Ridge, TN). This detector was connected to (1) regular detector electronics setup using a pulse-shaping amplifier (0.5  $\mu$ s shaping time) and a multi-channel analyzer, and (2) to the pile-up prevention circuit ("PPC") of the present invention.

The circuit may be interfaced to a computer. For example, the present invention was connected to a PC computer with a 133 MHz PENTIUM processor and a high-speed input/output board for the studies. The data acquisition software was written with LABVIEW (National Instruments, Austin, Texas). A  $^{99m}\text{Tc}$  point source (140 KeV gamma ray) in air was used for all the count-rate studies. The pile-up-prevention circuit of the present invention may also be used to measure the number of nonpile-ups, single pile-ups, and multiple pile-ups.

For comparison studies, circuits implementing three methods were setup and tested. The methods included: (i) the pulse-shaping method (500ns); (ii) delay-line clipping (256 ns); and (iii) the method according to the present invention. The pulse-height spectra, energy resolution and acquired true-count fractions were collected as a function of true-even rates using a  $^{99m}\text{Tc}$  (140 KeV) point source, which decays with a 6 h half-life. This half-life was used to calculate the true-even rate impinging the detector when very high count-rate data were acquired.

### EXAMPLE I

The pulse height spectra of  $^{99m}\text{Tc}$  (140 KeV gamma-ray) for the pile-up prevention circuit of the present invention and the pulse-shaping method (500ns) are shown in FIGS. 12A and 12B. FIG. 12A show results of the pulse-shaping method, and FIG. 12B show results of the present invention. At 50 Kcps, both measured 10.9% energy resolution. At 120 Kcps, the pulse-shaping method (FIG. 12A) started to demonstrate pile-up at the higher energy side and distortion at the lower energy side. Above 200-300 Kcps, the spectra were not usable. The PPC method spectra (FIG. 12B) maintained the spectra shape even at 2 Mcps. This first study demonstrated that the method of the present invention can significantly extend scintillation counting rates. However, the energy resolution was poor at 2 Mcps for this first study. Part of the reason was due to DC-level instabilities because the resistor divider of the photomultiplier (PMT) in the ORTEC probe was AC-coupled and the biasing-current was not designed for such a high count rate (there was no reason to design for such high rates previously because of NaI(Tl) pulse pile-up). Another reason was that the prototype board incorporating the apparatus of the present invention was built on a simple board with holes and the wiring was done with wirewrapping wire. A third reason was that a single integrator was used, which may not have been sufficiently discharged if the last event had a very high energy. Both the PMT divider and the prototype circuit were subsequently improved. Dual integrators and ADC channels ("ping-pong") were provided to assure complete discharge of the integrators. The pulse height spectra were measured again for the method of the present invention, and the energy resolution was significantly

improved, as shown in FIGS. 13A-D. Even at 2 Mcps (FIG. 13C), the energy resolution achieved 15% compared to 10.9% at 50 Kcps; part of the reason for this difference is the shorter average signal collection time when the count rate is very high. At lower count rates, the integration time for most pulses were integrated to 1 $\mu$ s, while at 2 Mcps, the average time-lapse between events was 500 ns which implies that the average signal collection time would be equal to or less than 500 ns.

## EXAMPLE II

In a second study, the method of the present invention was compared with a delay-line clipping (DLC) experimental setup (using NIM electronic modules and a 256 ns clipping). The energy resolution results are shown in FIG. 14. At lower count rates, the PPC method had better resolution than the DLC method, while at 2 Mcps the energy resolution was about the same. The study also generated the photopeak-count rates, *i.e.* the non-pile-up good counts acquired by each method using a narrow energy window, as a function of true-count rates, as shown in FIG. 15. For the method of the present invention at 2 Mcps, the photopeak counts were 85% of the total incident counts, while the DLC had only 25% good counts. The low photopeak fraction in the DLC method was due to a high 75% count-loss due to pile-ups which placed the pile-up-event energy above the photopeak acceptance-window. The DLC method delayed the occurrence of pile-up due to its shorter pulse width, but at very high count rates, its fixed pulse-width still made it more vulnerable to pile-ups (FIG. 16) than the PPC method, which has a variable signal-integrating time and count-recovery capability. Hence, the photopeak counts collected by the PPC method is 3.3 times more than the DLC method at 2 Mcps. The pulse-shaping method only maintained 35% at 0.25 Mcps, and there was no useful photopeak data above 0.25 Mcps because of pile-ups (FIG. 15).

The importance of an effective remnant light subtraction for every triggering event is demonstrated in FIGS. 17A-F, especially at very high counting rates. At 600 Kcps, a simple weighted-sum method (SWS) without remnant subtraction yielded an acceptable but worse spectrum compared to the present invention. However, at high counting rates 1-3.5 Mcps, the energy measurement errors of SWS are significant. Above

2 Mcps for NaI(Tl), most events are riding on the signal of other events (see Table 2 and FIG. 16), and SWS would measure large energy errors on most events and cause severe count-loss outside the photopeak acceptance window. The SWS energy measurement errors demonstrated here (shaded spectra in FIGS. 17A-F) should be similar to that shown in U.S. Patent No. 5,430,406 discussed above. FIGS. 15-17 demonstrates the count and energy recovery/restoration capabilities of the present invention.

### EXAMPLE III

As shown in FIGS. 18A-C, a Monte Carlo simulation of the torso (with liver, kidneys, and bladder) was performed to demonstrate the effect of pile-ups in imaging at 2 Mcps for: (a) the method of the present invention (FIG. 18A); (b) the pulse-clipping method (256 ns) (FIG. 18B); and (c) the regular fixed integration method (1  $\mu$ s) (FIG. 18C). The study simulates the electronic processing of each method in a gamma camera. The results are shown in FIGS. 18A-C. At 2 Mcps, the present invention method showed all the uptake organs; the conventional method shows no organs; the pulse-clipping result was in between the two, with high pile-up artifacts between the two kidneys and diminishing activities in the bladder and kidneys.

Besides preventing measurement errors from pile-ups, the present invention allows data acquisition time to be shortened significantly. This fast data acquisition is important in commercial/medical applications, as data acquisition time can be significantly reduced. As discussed above, the count loss with the conventional method is very high above 160,000 cps and saturates at 200,000 cps. Although the data appears to indicate that the conventional method can approach 200,000 cps, the pile-up fraction at this limit is very large, and the photopeak is badly distorted. Therefore, a large percentage of the data is not useful. The present invention can count over 2,000,000 cps with very little count loss, while preserving a very good photopeak, and all of the data is useful. Thus, the present invention is a significant improvement over the prior art.

Data from an embodiment of the present invention demonstrate that the invention significantly improves measurement of the energy of gamma rays at very high

count-rates, even when multiple pile-up approaches 60% and single pile-up approaches 17% (nonpile-up at only 13%) at 2,000,000 cps. The data indicate that a 10-20 fold improvement in count-rate capability is feasible with the present invention.

These studies discussed above showed that the methods of the present invention are feasible. Compared to the pulse-shaping method, it has a 15 times higher counting rates, and yet at low counting rates, the PPC energy resolution is the same as the standard pulse-shaping method. The PPC energy resolution only degraded slowly from 10.9% at 50 Kcps to 15% at 2000 Kcps. Compared to the pulse-clipping technique, the PPC method has better energy resolution at regular count rates (10.9% vs. 15%) and 3.3 times higher photopeak (non-pile-up) counts at 2000 Kcps. The count and energy recovery capability of the present invention at very high count rates is important for extending the maximum count rates of scintillation detectors.

All of the methods and apparatus disclosed and claimed herein can be made and executed without undue experimentation in light of the present disclosure. While the compositions and methods of this invention have been described in terms of preferred embodiments, it will be apparent to those of skill in the art that variations may be applied to the methods and apparatus and in the steps or in the sequence of steps of the method described herein without departing from the concept, spirit and scope of the invention. All such similar substitutes and modifications apparent to those skilled in the art are deemed to be within the spirit, scope and concept of the invention as defined by the appended claims.

## APPENDIX

### (1) Energy Pile-up Prevention Formulae Algebra

Suppose at time  $t_j$ , a  $j$ -th gamma ray ( $\gamma_j$ ) is detected. The  $j$ -th gamma ray is a pile-up event on top of two preceding gamma rays  $\gamma_{j-1}$  and  $\gamma_{j-2}$ . From *Knoll*, 1979, the instantaneous scintillation emission of a scintillator after detecting a gamma ray is

$$\xi(t) = \frac{E}{\tau} e^{-(t-t_j)/\tau} \quad (1)$$

where E is the total scintillation signal generated by the scintillator, and E is proportional to the energy of the gamma ray detected. The instantaneous signal at time  $t_j$  for  $t_j < t < t_{j+1}$ , contains emission from three gamma rays  $\gamma_{j-2}$ ,  $\gamma_{j-1}$ ,  $\gamma_j$ , and is given by,

$$\begin{aligned} q(t) &= \frac{E_j}{\tau} e^{-(t-t_j)/\tau} + \frac{E_{j-1}}{\tau} e^{-(t-t_{j-1})/\tau} + \frac{E_{j-2}}{\tau} e^{-(t-t_{j-2})/\tau} \\ &= \frac{E_j}{\tau} e^{-(t-t_j)/\tau} + \left[ \frac{E_{j-1}}{\tau} e^{-(t-t_{j-1})/\tau} \right] e^{-(t-t_j)/\tau} + \left[ \frac{E_{j-2}}{\tau} e^{-(t-t_{j-2})/\tau} \right] e^{-(t-t_j)/\tau} \end{aligned} \quad (2)$$

$$5 \quad = \left[ E_j + E_{j-1} e^{-(t_j-t_{j-1})/\tau} + E_{j-2} e^{-(t_j-t_{j-2})/\tau} \right] \left[ \frac{1}{\tau} e^{-(t-t_j)/\tau} \right] \quad (3)$$

If  $q(t)$  is amplified or weighted by an amplification factor  $\tau$ :

$$\tau q(t) = \left[ E_j + E_{j-1} e^{-(t_j-t_{j-1})/\tau} + E_{j-2} e^{-(t_j-t_{j-2})/\tau} \right] \left[ e^{-(t-t_j)/\tau} \right] \quad (4)$$

If  $q(t)$  is only integrated from  $t_j$  to  $t$  (this implicitly means that the integrator was discharged or zeroed just before  $t_j$ ) from equation (3),

$$10 \quad \int_{t_j}^t q(u) du = \left[ E_j + E_{j-1} e^{-(t_j-t_{j-1})/\tau} + E_{j-2} e^{-(t_j-t_{j-2})/\tau} \right] \left[ 1 - e^{-(t-t_j)/\tau} \right] \quad (5)$$

If equation (4) and (5) are summed, the weighted sum within the time period of gamma-ray  $\gamma_j$  will be:

$$S_j \equiv \tau q(t) + \int_{t_j}^t q(u) du = \left[ E_j + E_{j-1} e^{-(t_j-t_{j-1})/\tau} + E_{j-2} e^{-(t_j-t_{j-2})/\tau} \right] \left[ e^{-(t-t_j)/\tau} + 1 - e^{-(t-t_j)/\tau} \right] \quad (6)$$

15 The same analysis is now applied to the preceding time period where  $t_{j-1} < t < t_j$  (the time period of  $\gamma_{j-1}$ ). From equation (1) and FIG. 13, the instantaneous signal after the onset of  $\gamma_{j-1}$  is:

$$q(t) = \frac{E_{j-1}}{\tau} e^{-(t-t_{j-1})/\tau} + \frac{E_{j-2}}{\tau} e^{-(t-t_{j-2})/\tau} \quad (7)$$

The weighted sum for this period is:

$$S_{j-1} = \tau q(t) + \int_{j-1}^t q(u) du \quad (8)$$

Hence  $S_{j-1}$  is the term inside the square bracket of equation (6), and equation (6) becomes:

$$5 \quad S_j = E_j + S_{j-1} e^{-(t_j - t_{j-1})/\tau}$$

or 
$$\boxed{E_j = S_j - S_{j-1} e^{-(t_j - t_{j-1})/\tau}} \quad (9)$$

Hence the energy deposition of the  $j$ -th gamma ray can be derived from the  $j$ -th weighted sum minus the preceding weighted sum decreased by a emission decay factor, provided that the integrator is discharged immediately before the arrival of the  $j$ -th gamma. Note that the weighted-sum  $S_j$  as shown in equation (6) is not a measurement of energy of the  $j$ -th event, it is the measurement of the total energy trapped inside the scintillator just after the  $j$ -th event hit; *i.e.*,  $S_j$  includes the energy of the  $j$ -th event and all the remnant energy of previous multiple-pile-ups when the  $j$ -th event is deposited.

Equation (9) is an exact solution, regardless of how many previous gamma tails on which the  $j$ -th gamma is riding. The present invention thus provides the first method that attempts to recover **all** triggering events and compensating for multiple-pile-up exactly. Thus, the present invention may achieve higher count rates than possible with known methods.

Equation (9) can be easily understood intuitively. If there is no pile-up,  $S_j$  is the detected energy of the  $j$ -th gamma ray (independent to when  $S_j$  is sampled). If there are pile-ups,  $S_j$  is the **total energy** in the detector in the  $j$ -th period (independent to when  $S_j$  is sampled), which includes the detected energy of the  $j$ -th gamma ray and the remnant emission of **all** the preceding gamma ray. This  $j$ -th remnant emission is simply the total



energy in the preceding period ( $S_{j-1}$ ) decreased by the emission decay factor governed by the time difference between the  $j$ -th gamma and its preceding gamma.

## (2) Pile-up Prevention Formulae for Gamma-Camera Anger-Positioning

5 In Anger-position method, assuming that no pile-up occurs, the  $\gamma$ -location is estimated by the centroid calculation:

$$X = \frac{\sum_i X_i E_i}{\sum_i E_i} = \frac{X'}{Z} \quad ; \quad Z \equiv \sum_i E_i \quad (10)$$

where  $E_i$  is the total signal received by Photomultiplier- $i$  (PMT- $i$ ) when a  $\gamma$ -ray is detected, and  $X_i$  is the physical location of PMT- $i$ .  $X'$  is the prenormalized pseudo-position signal which needs to be normalized by the total detected energy  $Z$  of the  $\gamma$ -ray, to generate the centroid location  $X$ .  $X'$  can be defined as:

$$X' \equiv X_+ - X_-$$

If at time  $t_1$ , the first  $\gamma$ -ray with energy  $Z_1$  is detected, the total energy (after signal integration) distributed to PMT-1, PMT-2,  $\dots$ , PMT- $n$  are  $E_{11}$ ,  $E_{12}$ ,  $\dots$ ,  $E_{1n}$ , respectively. The total integrated energy detected by the camera is:

$$15 \quad Z_1 = \sum_i E_{1i} \quad (11)$$

From equation (10), the prenormalized position  $X'$  for conventional Anger positioning is:

$$X'_1 = \sum_i X_i E_{1i} \quad (12)$$

In conventional Anger-position,  $E_{Ii}$  in equation (12) are the integrated PMT signals (integrated for 0.6-1  $\mu$ s), and therefore  $X'_1$  in equation (12) is a static quantity obtained after a fixed 1  $\mu$ s integration time. Hence, if there is a pile-up during the integration time,  $X'_1$  will be wrong. In the method of the present invention, the  
 5 instantaneous prenormalized position  $X'(t)$  is used:

$$X'_1(t) = \sum_i X_i \frac{E_{Ii}}{\tau} e^{-(t-t_i)/\tau}$$

The weighted instantaneous prenormalized position is:

$$\tau X'_1(t) = \sum_i X_i E_{Ii} e^{-(t-t_i)/\tau} \quad (13)$$

The integrated prenormalized position is:

$$10 \quad \int_1 X'_1(u) du = \sum_i X_i E_{Ii} (1 - e^{-(t-t_i)/\tau}) \quad (14)$$

If equation (13) and (14) are summed, the weighted Sum of  $X'_1(t)$  is given by:

$$SX'_1 = \tau X'_1(t) + \int_1 X'_1(u) du = \sum_i X_i E_{Ii} \quad (15)$$

From equation (12) and equation (15), we obtain  $SX'_1 = X'_1$ . Hence, if the event  $\gamma_1$  is not a pile-up on previous events, the weighted sum,  $SX'_1$ , is also an estimation of the  
 15 Anger centroid-location for the gamma ray ( $\gamma_1$ ), just like static quantity  $X'_1$  in equation (12) for conventional Anger-method. Furthermore, the weighted sum  $SX'_1$  is a time independent quantity which can be sampled at anytime and still provide an unbiased estimation of the prenormalized location. Therefore sum position  $SX'_1$  is just like its energy-counterpart weighted sum  $S_i$  in equation (6).

20 However, if the event ( $\gamma_1$ ) is a pile-up on a previous gamma-ray ( $\gamma_0$ ), the prenormalized position weighted sum will not be equal to the true prenormalized position

of  $\gamma_1$ , because  $SX'_1$  contains the remnant prenormalized position weighted-sum  $SX'_0$  generated by the remnant scintillation of  $\gamma_0$ . The  $\gamma_0$  remnant position weighted-sum (a position vector) adds an error-position vector to the  $\gamma_1$  position. The magnitude of this error vector increases if the time lapse ( $t_1-t_0$ ) between  $\gamma_0$  and  $\gamma_1$  decreases. These

5 positioning errors are especially severe when the count rates are very high that (a) most of the events ride on the remnant of previous events and (b) multiple-event pile-ups are the norm. Hence, the general solution to derive the true prenormalized position of each event  $\gamma_1$  requires the subtraction of the remnant-position signals of previous events from the prenormalized position weighted-sum,  $SX'_i$  and  $SY'_i$ , triggered by each  $\gamma_i$ .

10 In the general case where any event can be riding on the signals of one or more previous events, the weighted instantaneous prenormalized position signal triggered by an incident event  $\gamma_m$  is given by the superposition,

$$\tau X'_m(t) = \sum_i X_i E_{m,i} e^{-(t-t_m)/\tau} + \sum_i X_i E_{m-1,i} e^{-(t-t_{m-1})/\tau} + \sum_i X_i E_{m-2,i} e^{-(t-t_{m-2})/\tau} + \dots \quad (16)$$

where  $t_m < t < t_{m+1}$ . The corresponding integrated prenormalized position signal is given

15 by:

$$\begin{aligned} \int_{t_m}^t X'_m(u) du = & \sum_i X_i E_{m,i} (1 - e^{-(t-t_m)/\tau}) \\ & + \sum_i X_i E_{m-1,i} e^{-(t_m-t_{m-1})/\tau} (1 - e^{-(t-t_m)/\tau}) \\ & + \sum_i X_i E_{m-2,i} e^{-(t_m-t_{m-2})/\tau} (1 - e^{-(t-t_m)/\tau}) \\ & + \dots \end{aligned} \quad (17)$$

This integration can be carried out by first zeroing the integrator just before the onset of  $\gamma_m$ , and then restarting the integration immediately.

Summing equations (16) and (17), the weighted sum of the prenormalized

20 position signal is:

$$SX'_m = \sum_i X_i E_{m,i} + \sum_i X_i E_{m-1,i} e^{-(t_m - t_{m-1})/\tau} + \sum_i X_i E_{m-2,i} e^{-(t_m - t_{m-2})/\tau} + \dots \quad (18)$$

$$= \sum_i X_i E_{m,i} + e^{-(t_m - t_{m-1})/\tau} \left[ \sum_i X_i E_{m-1,i} + \sum_i X_i E_{m-2,i} e^{-(t_{m-1} - t_{m-2})/\tau} + \dots \right] \quad (19)$$

It can be observed that the term inside the bracket in equation (19) is simply  $SX'_{m-1}$  (see equation (18)). The recurring relationship reduced equations (19) and (12) to

$$\begin{aligned} 5 \quad SX'_m &= \sum_i X_i E_{m,i} + e^{-(t_m - t_{m-1})/\tau} SX'_{m-1} \\ &= X'_m + e^{-(t_m - t_{m-1})/\tau} SX'_{m-1} \end{aligned}$$

Therefore, the prenormalized Anger-logic position  $(X'_m, Y'_m)$  of  $\gamma_m$  can be derived from the weighted prenormalized position sums of  $\gamma_m$ ,  $\gamma_{m-1}$ , and the time-lapse between the two events,

$$\begin{aligned} X'_m &= SX'_m - SX'_{m-1} e^{-(t_m - t_{m-1})/\tau} \\ Y'_m &= SY'_m - SY'_{m-1} e^{-(t_m - t_{m-1})/\tau} \end{aligned} \quad (20)$$

10 Since the derivation of equation (20) assumes that  $\gamma_m$  is riding on any numbers of preceding events, equation (20) provides an exact determination of the position of  $\gamma_m$ , even when it is part of a multiple-event pile-up.

The pile-up-free prenormalized position  $(X'_m, Y'_m)$  will then be renormalized by the energy of  $\gamma_m$  as in conventional Anger-logic to remove the energy-scaling effect in  
15 generating the position of  $\gamma_m$ . Hence, the energy of  $\gamma_m$  also has to be extracted correctly from a multiple-event pile-up if the position of  $\gamma_m$  is to be decoded accurately. The pile-up-free energy of  $\gamma_m$  can also be derived from the energy weighted-sum of  $\gamma_m$ ,  $\gamma_{m-1}$ , and the time lapse between these two events,

$$E_m = S_m - S_{m-1} e^{-(t_m - t_{m-1})/\tau} \quad (21)$$

where  $S_m$  is the sum of the weighted instantaneous scintillation signal and the integral of the instantaneous scintillation signal.

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